Elementary Particle Physics: Assignment # 2Due TUESDAY Feb 9th at 11:30 pm (before starting class)

(1) The momentum expansion of a free scalar complex field

$$\Phi(x) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2Ep}} \left[a_p e^{-ipx} + b_p^{s\dagger} e^{ipx} \right] \text{ with } E_p = \sqrt{|\vec{p}|^2 + m^2}$$

 a_p and b_p verify: $[a_p, a_q^{\dagger}] = [b_p, b_q^{\dagger}] = (2\pi)^3 \delta^3(p-q)$ and all other commutators vanish. The propagator is defined

$$i\Delta_F(x-y) \equiv \langle 0|T\left[\Phi(x)\Phi^*(y)\right]|0\rangle \equiv \Theta(x_0-y_0)\langle 0|\Phi(x)\Phi^*(y)|0\rangle + \Theta(y_0-x_0)\langle 0|\Phi(y)^*\Phi(x)|0\rangle$$

Using the expansion of the fields and the commutation relations above show that

$$\Delta_F(x-y) = \Theta(x_0 - y_0) \Delta_+(x-y) - \Theta(y_0 - x_0) \Delta_-(x-y)$$

with

$$\Delta_{\pm}(x-y) = \mp i \int \frac{d^3p}{(2\pi)^3} \frac{1}{2E_p} e^{\mp ip(x-y)}$$

where Δ_+ comes from the term $\langle 0|a_pa_p^{\dagger}|0\rangle$ and Δ_- comes from the term $\langle 0|b_pb_p^{\dagger}|0\rangle$

(2) In the chiral representation the 4-spinors for a fermion with momentum $\vec{p} = |\vec{p}|(\sin\theta\cos\phi, \sin\theta\sin\phi, \cos\theta)$ with positive and negative helicity are:

$$u^{1,2}(\vec{p}) = u^{\pm}(\vec{p}) = \begin{pmatrix} \sqrt{E \mp |\vec{p}|} \ \xi_p^{\pm} \\ \sqrt{E \pm |\vec{p}|} \ \xi_p^{\pm} \end{pmatrix} \quad \text{with} \quad \xi_p^{+} = \begin{pmatrix} \cos\frac{\theta}{2} \\ e^{i\phi} \sin\frac{\theta}{2} \end{pmatrix} \quad \xi_p^{-} = \begin{pmatrix} -e^{-i\phi} \sin\frac{\theta}{2} \\ \cos\frac{\theta}{2} \end{pmatrix}$$

and the corresponding 4-spinors for the anti-fermion are: $v^{1,2}(\vec{p}) = v^{\pm}(\vec{p}) = \pm \begin{pmatrix} \sqrt{E \pm |\vec{p}|} \xi_p^{\mp} \\ -\sqrt{E \mp |\vec{p}|} \xi_p^{\mp} \end{pmatrix}$ Using these expressions evaluate by direct calculation

 $\overline{u}^{s}(\vec{p})v^{r}(-\vec{p})$ for the four helicity combinations

(3) The Lagrangian and momentum expansion for the free Dirac field are

$$\mathcal{L} = N\left[\bar{\psi}(x)\left(i\gamma^{\mu}\partial_{\mu}-m\right)\psi(x)\right] \qquad \psi(x) = \int \frac{d^{3}p}{(2\pi)^{3}} \frac{1}{\sqrt{2Ep}} \sum_{s} \left[u^{s}(p)c_{p}^{s}e^{-ipx} + v^{s}(p)d_{p}^{s\dagger}e^{ipx}\right]$$

where c_p^s and d_p^s verify: $\left\{c_p^s, c_q^{r\dagger}\right\} = \left\{d_p^s, d_q^{r\dagger}\right\} = (2\pi)^3 \delta^3 (p-q) \delta^{rs}$ and all other anticommutators vanish.

4.1) Derive, starting from the Lagrangian above, the expression of the Hamiltionian and 3-momentum operators in terms of the fields ψ and $\overline{\psi}$.

4.2) Use the expansion of ψ above (and the corresponding for $\bar{\psi}$) and the anticommutation relations of d's and c's to show that

$$H = \int \frac{d^3p}{(2\pi)^3} \sum_s E_p \left[c_p^{s\dagger} c_p^s + d_p^{s\dagger} d_d^s \right]$$

4.3) What would you get if the fields c_p^s and d_p^s verified commutation (instead of anticommutation) relations?