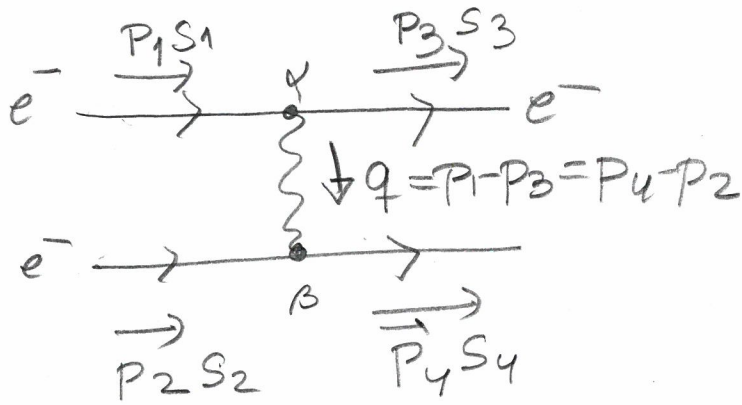
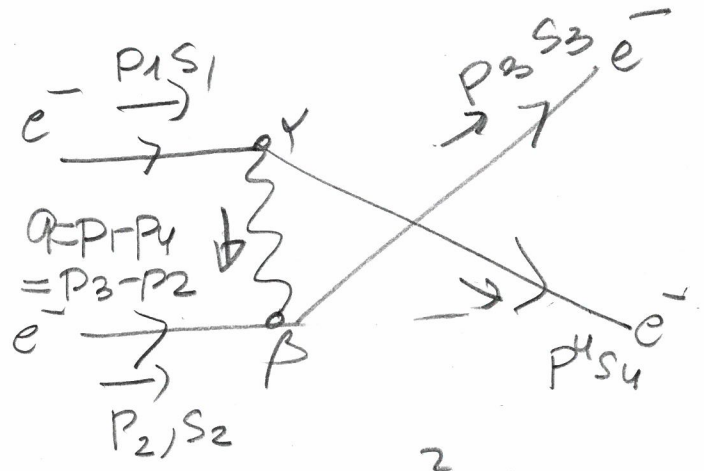


①  $e^- e^- \rightarrow e^- e^-$



(a)  $\Rightarrow q^2 = t$



(b)  $\Rightarrow q^2 = u$

$iM_a = \frac{e^2}{t} \bar{u}_3 \gamma^\alpha u_1 \bar{u}_4 \gamma_\alpha u_2$

48

exchanged identical fermions

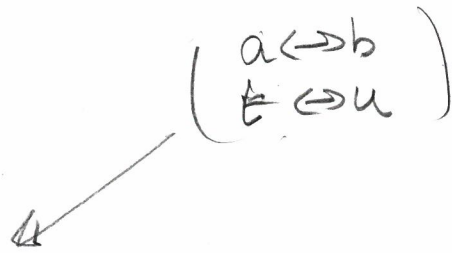
$iM_b = (-) \frac{e^2}{u} \bar{u}_3 \gamma^\alpha u_2 \bar{u}_4 \gamma_\alpha u_1$

$|\bar{M}|^2 = \frac{1}{4} \sum_{s_1 s_2 s_3 s_4} |M_a + M_b|^2 = |\bar{M}_a|^2 + |\bar{M}_b|^2 - 2 \text{Re} \bar{M}_a \bar{M}_b^*$

$|\bar{M}_a|^2 = \frac{e^4}{4t^2} \text{Tr}[\not{p}_3 \gamma^\alpha \not{p}_1 \gamma_\alpha] \text{Tr}[\not{p}_4 \gamma_\alpha \not{p}_2 \gamma_\alpha]$

$= \frac{2e^4}{t^2} (u^2 + s^2)$

$|\bar{M}_b|^2 = \frac{2e^4}{u^2} (t^2 + s^2)$



the inlets

$$\frac{-8s^2}{11}$$

$$-2\text{Re}(\overline{M}_a \overline{M}_b^*) = -\frac{2e^4}{4tu} \text{Tr} \left[ \begin{matrix} \rho_3 \gamma^0 \rho_2 \gamma^0 \rho_4 \gamma^0 \rho_2 \gamma^0 \end{matrix} \right]$$

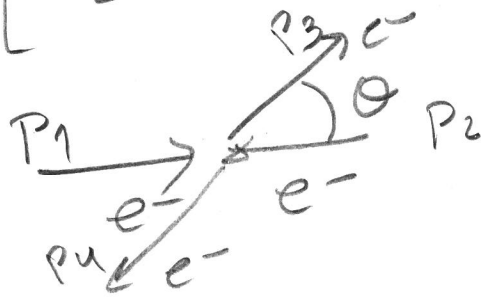
$$= + \frac{4s^2}{tu}$$

2 identical particles final state

So  $\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2 s} \frac{1}{2} |\overline{M}|^2$

$$= \frac{e^4}{64\pi^2 s} \left[ \frac{s^2 + u^2}{t^2} + \frac{s^2 + t^2}{u^2} + \frac{2s^2}{tu} \right]$$

In COM



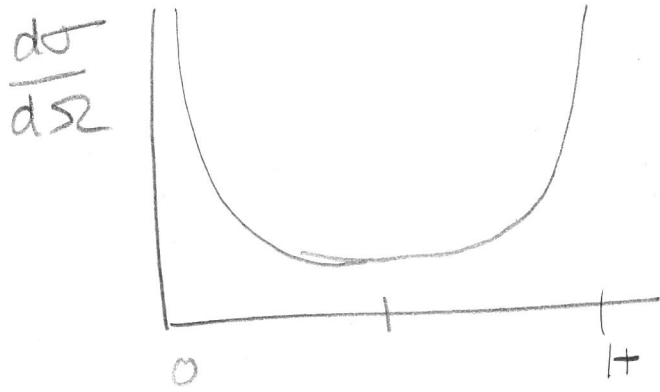
$$t = -\frac{s}{2} (1 - \cos\theta) \quad \text{and} \quad u = -\frac{s}{2} (1 + \cos\theta) = -s \cos^2 \frac{\theta}{2}$$

$$= -s \sin^2 \frac{\theta}{2}$$

$$\frac{d\sigma}{d\Omega} \Big|_{\text{COM}} = \frac{e^4}{64\pi^2 s} \left[ \frac{1 + \sin^4 \frac{\theta}{2}}{\cos^4 \frac{\theta}{2}} + \frac{1 + \cos^4 \frac{\theta}{2}}{\sin^4 \frac{\theta}{2}} + \frac{1}{\sin^2 \frac{\theta}{2} \cos^2 \frac{\theta}{2}} \right]$$

F(theta)

② This CS becomes  $\infty$  at  $\theta=0, \pi$  so  
 and it must be symmetric  
 when  $\theta \leftrightarrow \pi - \theta$   
 since  $e^- \leftrightarrow e^-$



identical

numerically using  $\alpha = \frac{e^4}{4\pi} \approx \frac{1}{137}$   $S = (50 + 50)^2 = (100 \text{ GeV})^2$   
 $\Delta \text{ in NU}$

$$\frac{e^4}{64\pi^2 S} = \frac{\alpha^2}{4S} = 13.3 \times 10^{-10} \text{ GeV}^{-2} \times (197 \text{ MeV fm})^2$$

$$= 5.16 \times 10^{-11} \text{ fm}^2 = 5.16 \times 10^{-41} \text{ m}^2$$

$$= 0.516 \text{ pb}$$

$$F(\theta) = \begin{cases} \theta=0 & \infty \\ \theta=10 & 3.5 \times 10^4 \\ \theta=45 & 98 \\ \theta=90 & 18 \\ \theta=137 & 98 \\ \theta=170 & 3.5 \times 10^4 \\ \theta=180 & \infty \end{cases}$$

(3)

$$\frac{N_{e^-}(\theta)}{\text{sr min}} = \alpha \times \frac{d\sigma}{d\Omega}(\theta) \times \text{time}$$

$$\alpha = \frac{N^2 \times f}{A} = \frac{(10^{11})^2}{400 \times (10^{-6} \text{ m})^2} \times \frac{1}{25 \times 10^{-9} \text{ s}} = 10^{39} \text{ m}^{-2} \text{ s}^{-1}$$

2 per collision →

$$\frac{N_{e^-}(\theta)}{\text{sr min}} = 2 \times 5.16 \times 10^{-4} \text{ m}^2 \times 10^{39} \text{ m}^{-2} \text{ s}^{-1} \times 60 \text{ s} F(\theta)$$

1 minute  
↓

$$= 6.2 \times F(\theta) = \begin{cases} 1.12 & \text{for } \theta = 90^\circ \\ 2.1 \times 10^9 & \text{for } \theta = 1^\circ \end{cases}$$

At  $\theta = 1$  deg we are close to  $\theta = 0$  for which the CS becomes  $\infty$  due to  $\infty$  range of em interaction