

PHY682 Special Topics in Solid-State Physics: Quantum Information Science

Lecture time: 2:40-4:00PM Monday & Wednesday

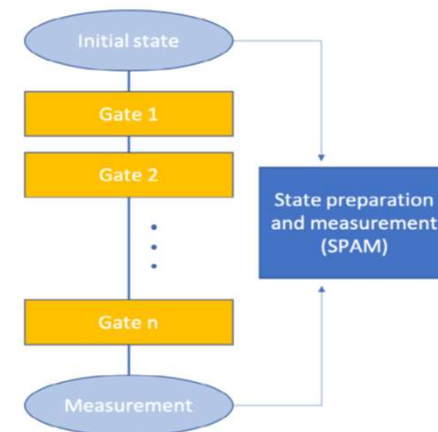
Today 11/9:

1. Final presentation discussion (within each group)
2. Finish Week 10's topics (today on tomographic tools)
3. Begin Week 11's topics (quantum phase estimation and applications)

Tomographic tools

- Crucial to ensure proper functioning of QC and correctness of results

- ❖ State preparation
⇒ State tomography
- ❖ Gate operations
⇒ Process tomography
- ❖ Measurement (i.e. detectors)
⇒ Detector tomography



- Note: detector tomography is often ignored, but important to extract correct computational outcomes

Quantum state tomography

□ Estimate unknown state (given multiple copies)

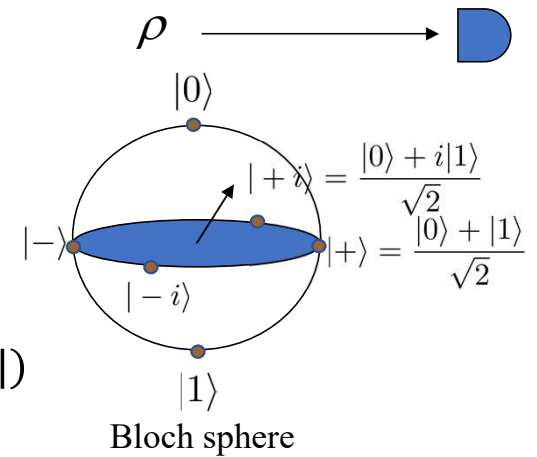
□ One qubit

$$\rho = \frac{1}{2}(I + r_x\sigma_x + r_y\sigma_y + r_z\sigma_z)$$

$$r_x = \text{tr}(\rho\sigma_x) = \text{tr}(\rho|+\rangle\langle+|) - \text{tr}(\rho|-\rangle\langle-|)$$

$$r_y = \text{tr}(\rho\sigma_y) = \text{tr}(\rho|+i\rangle\langle+i|) - \text{tr}(\rho|-i\rangle\langle-i|)$$

$$r_z = \text{tr}(\rho\sigma_z) = \text{tr}(\rho|0\rangle\langle 0|) - \text{tr}(\rho|1\rangle\langle 1|)$$



→ If one can measure the qubit in all three bases, can extract Bloch vector \vec{r}

□ Multi-qubits:

$$\rho_{2\text{-qubit}} = \frac{1}{4} \sum_{\mu\nu} r_{\mu\nu} \sigma_{\mu} \otimes \sigma_{\nu}, \quad r_{\mu\nu} = \text{tr}(\rho_{2\text{-qubit}} \sigma_{\mu} \otimes \sigma_{\nu})$$

→ Measure in product of bases (i.e. coincidence)

Quantum process tomography

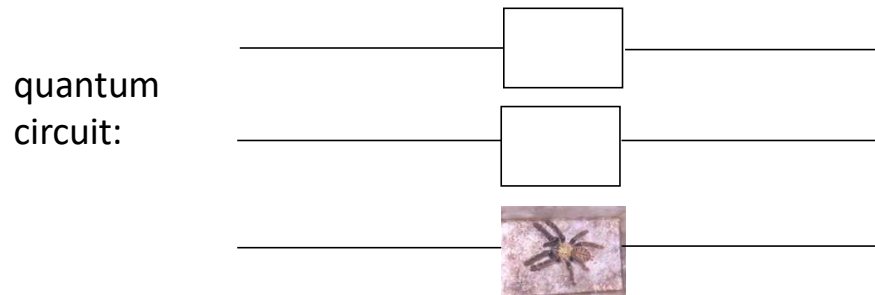
- Estimate unknown process (Black box)



$$E: \rho \rightarrow \sum_i E_i \rho E_i^\dagger$$

(quantum operations)

- Possible application: “debugging” quantum gates



Q: From measuring a limited number of different input states (but unlimited supply of each), is it possible to predict the result for a general input state?

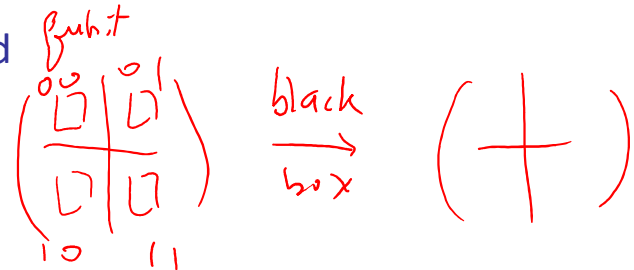
Three different ways of implementing
quantum process tomography (PT)

(I) Standard Quantum PT (SQPT)

➤ Idea: look at how each element gets transformed

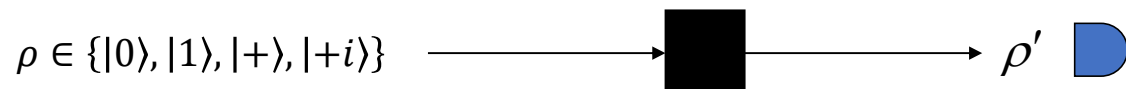
$$E: |j\rangle\langle k| \rightarrow \sum_i E_i |j\rangle\langle k| E_i^\dagger$$

$|j\rangle\langle k|$ for single qubit are



$$\begin{array}{l}
 |0\rangle\langle 0| \\
 |0\rangle\langle 1| = \underbrace{|+\rangle\langle +|}_{\downarrow} + \underbrace{i|+i\rangle\langle +i|}_{\downarrow} - \frac{1+i}{2} \underbrace{|0\rangle\langle 0|}_{\downarrow} - \frac{1+i}{2} \underbrace{|1\rangle\langle 1|}_{\downarrow} \\
 |1\rangle\langle 0| = \underbrace{|+\rangle\langle +|}_{\downarrow} - \underbrace{i|+i\rangle\langle +i|}_{\downarrow} - \frac{1-i}{2} \underbrace{|0\rangle\langle 0|}_{\downarrow} - \frac{1-i}{2} \underbrace{|1\rangle\langle 1|}_{\downarrow} \\
 |1\rangle\langle 1|
 \end{array}$$

➔ We only need four different inputs $|0\rangle, |1\rangle, |+\rangle, |+i\rangle$



to figure out the unknown action

(II) Entanglement assisted PT (EAAPT)

Q: Can we do better? Send in same state?

Yes, sending a Bell state!

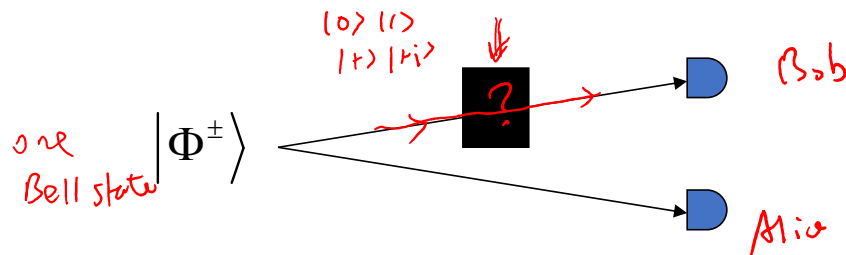
$$|\Phi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

$$\rho_{\Phi^+} = \frac{1}{4}(I \otimes I + \sigma_x \otimes \sigma_x - \sigma_y \otimes \sigma_y + \sigma_z \otimes \sigma_z)$$

$$\xrightarrow{E} \frac{1}{4}(\mathcal{E}(I) \otimes I + \mathcal{E}(\sigma_x) \otimes \sigma_x - \mathcal{E}(\sigma_y) \otimes \sigma_y + \mathcal{E}(\sigma_z) \otimes \sigma_z)$$

➤ By measuring output state, we can infer E on the complete set of matrices, thus on arbitrary input state

Alice Bob
 $|00\rangle + |11\rangle$
 to
 measure
 $\sigma_i \rightarrow \sigma_i$
 $+/- \rightarrow +/-$
 $\begin{pmatrix} + \\ - \end{pmatrix} / -i \rightarrow \begin{pmatrix} - \\ + \end{pmatrix} / +i$



Question: how can we use "remote-state preparation" to understand this?

(III) Ancilla assisted PT (AAPT)

■ Is entanglement necessary? $\rho_{\text{unentangled}}?$

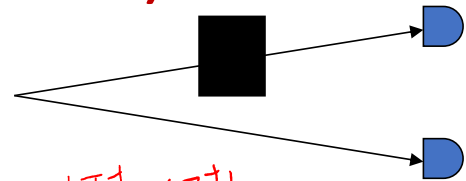
■ No, we can use an unentangled state!

➤ For example, use an unentangled Werner state

$$\rho_{\text{Werner}}(1/3) \equiv \frac{1}{6} I \otimes I + \frac{1}{3} \rho_{\Phi^+}$$

$$= \frac{1}{4} \left(I \otimes I + \frac{1}{3} \sigma_x \otimes \sigma_x - \frac{1}{3} \sigma_y \otimes \sigma_y + \frac{1}{3} \sigma_z \otimes \sigma_z \right)$$

$$\xrightarrow{E} \left(\frac{1}{4} \left(E(I) \otimes I + \frac{1}{3} E(\sigma_x) \otimes \sigma_x - \frac{1}{3} E(\sigma_y) \otimes \sigma_y + \frac{1}{3} E(\sigma_z) \otimes \sigma_z \right) \right)$$



$$\rho(\psi) = \gamma |\Phi^+\rangle \langle \Phi^+| + \frac{1-\gamma}{4} I$$

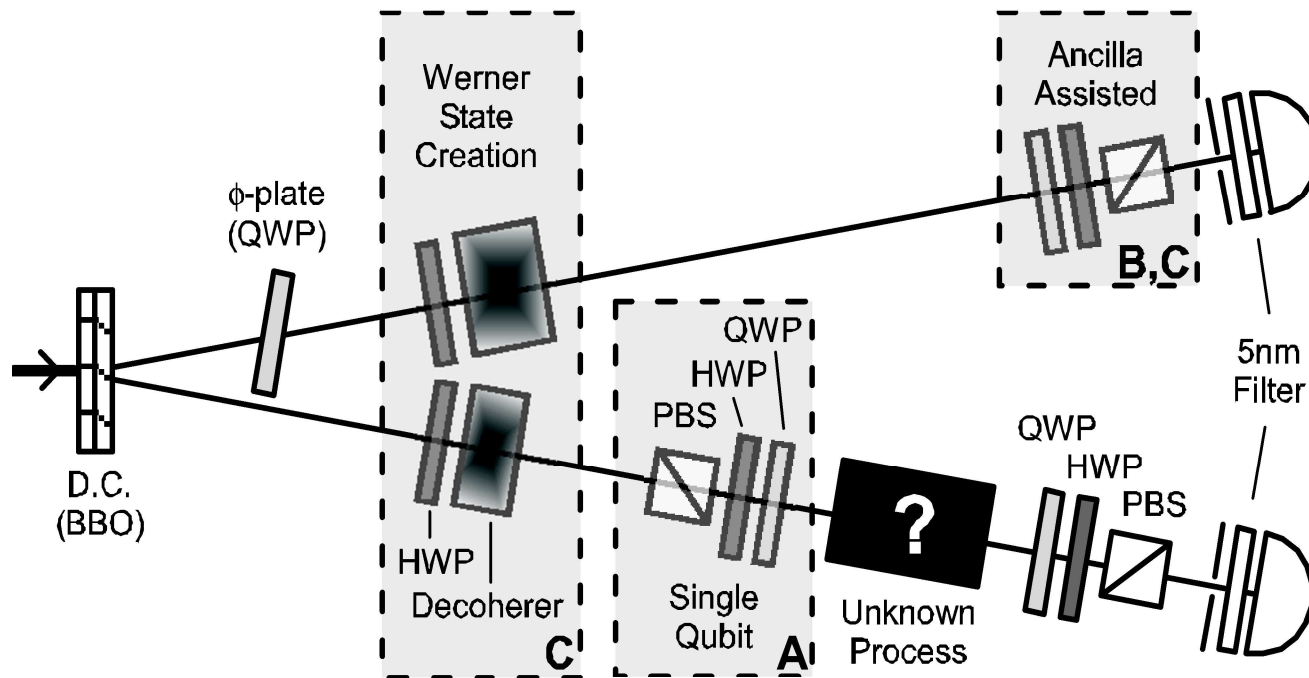
$\gamma \leq \frac{1}{3}$ not ent.

$$= \begin{pmatrix} \frac{1}{3} & 0 & 0 & \frac{1}{6} \\ 0 & \frac{1}{6} & 0 & 0 \\ 0 & 0 & \frac{1}{6} & 0 \\ \frac{1}{6} & 0 & 0 & \frac{1}{3} \end{pmatrix}$$

➤ Due to 1/3, the noise is expected to be higher than that of using a Bell state

Experimental setup for process tomography

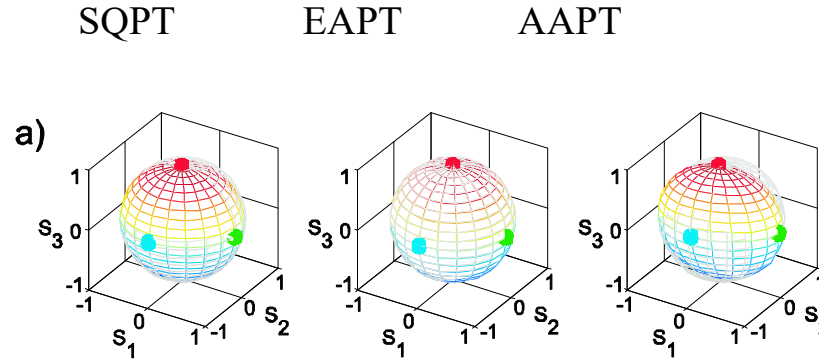
Altepeter, Branning, Jeffrey, Wei,
Kwiat, Thew, O'Brien, Nielsen, White,
Phys. Rev. Lett. 90, 193601 (2003)



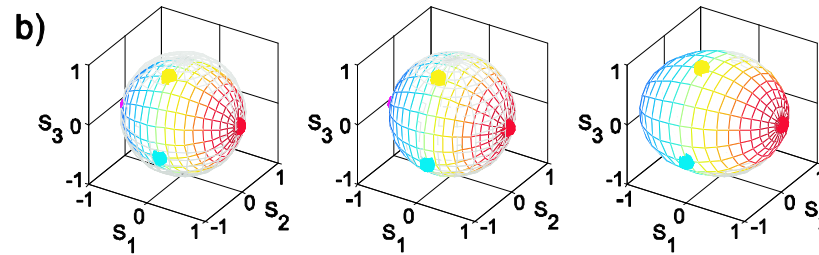
Experimental Results

[Altepeter et al., PRL'03]

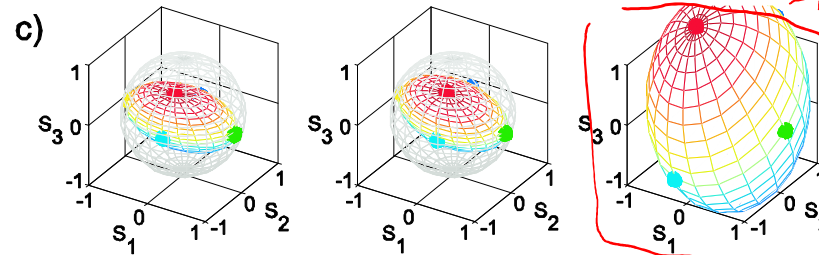
Identity



Quarter waveplate



Partial decoherence



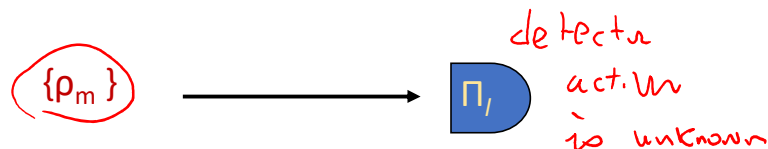
$\rightarrow r > 1$
 \Rightarrow not physical

$\mathcal{E}(P) = \sum_i E_i P E_i^\dagger$
 using the same degree of freedom to create Werner $\rho(\frac{1}{2})$

Detector tomography

[Fiurasek, PRA 64, 024102 (2001)]

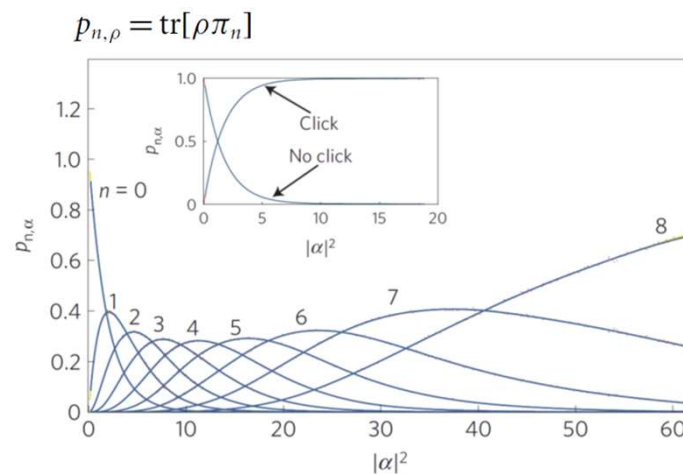
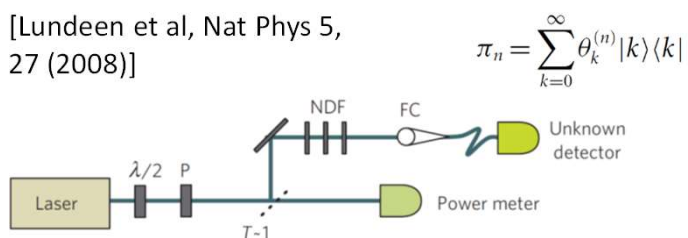
- Assumption: an informationally complete set of test states $\{\rho_m\}$ be prepared with error smaller than error in measurement



- Goal: infer when detector outcome l clicks, what is measured?

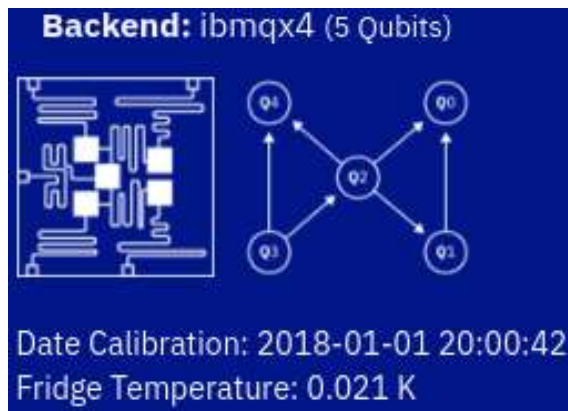
- Recent application in characterizing photon detectors

[Lundeen et al, Nat Phys 5, 27 (2008)]

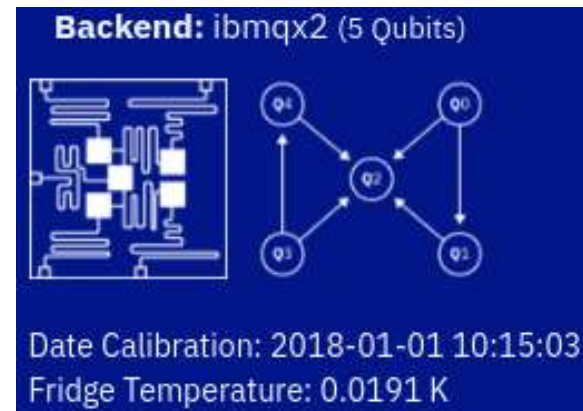


Carry out detector tomography on IBM Q Machines

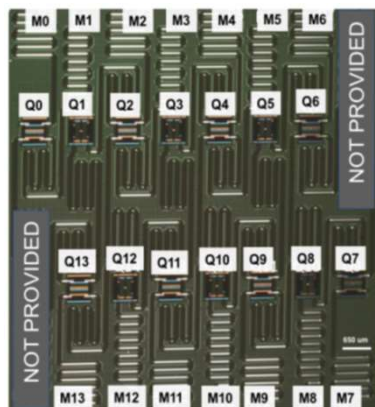
Tenerife



Yorktown



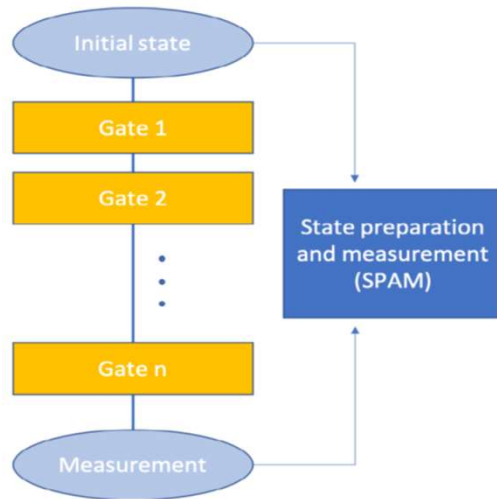
IBM Q 16 Melbourne (14 qubits)



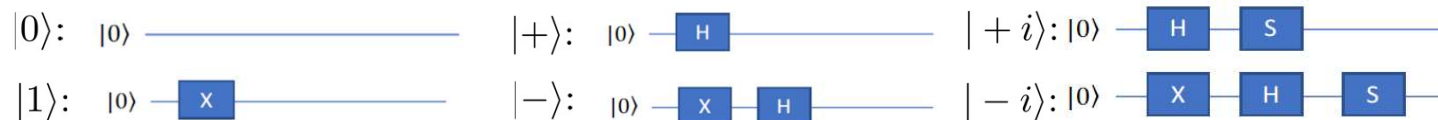
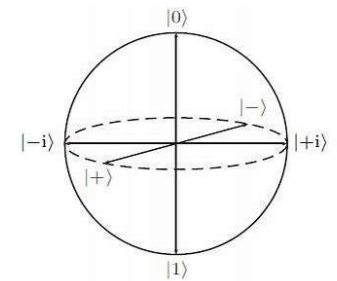
- Transmon qubits
(transmission line shunted plasma oscillation qubit)
- Coupled by resonators
- Qubit operations by microwave pulses; readout by measuring transmitted radiation

Detector tomography

[Fiurasek, PRA 64, 024102 (2001)]



- Assumption: an informationally complete set of test states $\{\rho_m\}$ be prepared with error smaller than error in measurement
- In IBM Q machines, measurement has large errors, but ground state $|0\rangle$ preparation and single-qubit gates are of high fidelity so can prepare $\{\rho_m = 0, 1, +, -, +i, -i\}$



- Goal: to characterize a 2-outcome positive operator valued measure (POVM): $\{\Pi_0, \Pi_1\}$

Maximum Likelihood for Detector Tomography*

[Fiurasek, PRA 64, 024102 (2001)]

- Prepare states ρ_m to characterize POVM elements $\pi_l \geq 0$

➤ Probability $p_{lm} = \text{Tr}(\rho_m \pi_l) \approx \frac{f_{lm}}{\sum_{l'} f_{l'm}}$ *from experimental data*

f_{lm} : number of times input ρ_m such that detector π_l clicks

- Use Maximize the likelihood function to iteratively find π_l

$$\mathcal{L}\{\Pi_l\} = \prod_l \prod_m (\text{Tr}[\Pi_l \rho_m])^{f_{lm}} \quad \text{constraint: } \sum_l \Pi_l = I$$

- Optimal π_l can be obtained by an iterative process:

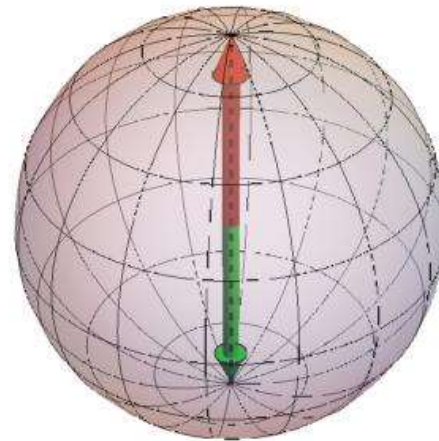
$$\Pi_l \leftarrow R_l \Pi_l R_l^\dagger \quad R_l = \hat{\lambda}^{-1} \sum_m \frac{f_{lm}}{p_{lm}} \rho_m$$
$$\hat{\lambda} = \left(\sum_{m',m,l} \frac{f_{lm} f_{lm'}}{p_{lm} p_{lm'}} \rho_{m'} \Pi_l \rho_m \right)^{1/2}$$

Perfect projectors onto 0 & 1

□ As a reminder, perfect 0 & 1 measurements are

$$\Pi_0 = |0\rangle\langle 0| = \frac{1}{2}I + \frac{1}{2}\sigma_z$$

$$\Pi_1 = |1\rangle\langle 1| = \frac{1}{2}I - \frac{1}{2}\sigma_z$$



Runs on IBMQ

- Prepare states that are \pm eigenstates of Pauli operators:

$$\{p_j = 0, 1, +, -, +i, -i\}$$

- Model of single-qubit detector (POVM):

[Chen, Farahzad, Yoo, & Wei, PRA 2019]

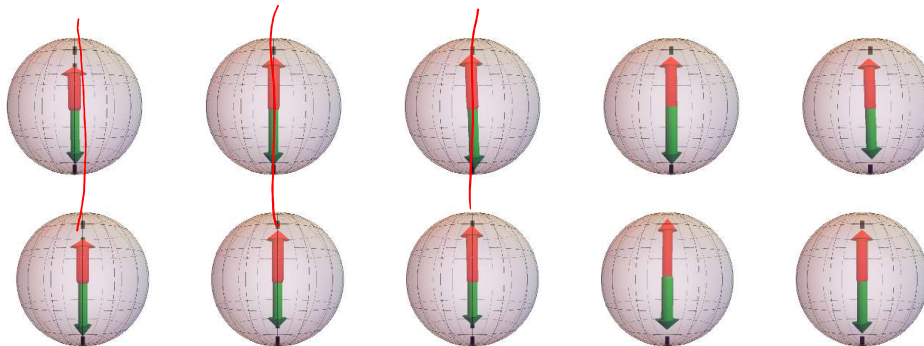
generalize
 $P : |0\rangle\langle 0| \xrightarrow{\text{detection}} M_e$
 $P_s = \langle 0|P|0\rangle \quad P_e = \text{Tr}(P M_e)$

$$|0\rangle\langle 0|$$

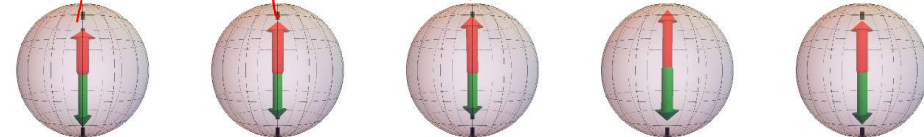
$$\Leftrightarrow \pi_0 = a_0^{(0)} I + \vec{a}^{(0)} \cdot \vec{\sigma}, \quad \pi_1 = a_0^{(1)} I + \vec{a}^{(1)} \cdot \vec{\sigma}$$

- Results on IBM QX4 (5 qubits):

➤ Measured individually



➤ measured 5 qubits in parallel



- ❖ Can be used to mitigate imperfect measurement (next slide)

Arrows indicate the vectors $\vec{a}^{(0)}$ or $\vec{a}^{(1)}$ averaged from 100 different runs; thickness of arrows represents magnitude of $a_0^{(0)}$ or $a_0^{(1)}$.

Mitigation of imperfect detectors ^{*}

- Assume detectors act independently (otherwise need multi-detector tomography)

[Chen, Farahzad, Yoo, & Wei, PRA 2019]

Distribution of measurement outcomes

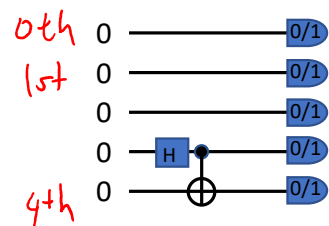
$$\begin{aligned}
 \tilde{P}_{(n_0, n_1, \dots, n_{N-1})} &\approx \text{Tr} \left[\rho \prod_{j=0}^{N-1} (a_{0,[j]} + a_{3,[j]}^{(n_j)} \sigma_{3,[j]}) \right] \\
 &= \sum_{(m_0, m_1, \dots, m_{N-1})} P_{(m_0, m_1, \dots, m_{N-1})} \prod_{j=0}^{N-1} (a_{0,[j]} + (-1)^{m_j} a_{3,[j]}^{(n_j)}), \\
 &= \sum_{\vec{m}} M_{\vec{n}; \vec{m}} P_{\vec{m}}
 \end{aligned}$$

measured $\tilde{P}_{(n_0, n_1, \dots, n_{N-1})}$ \approx $\text{Tr}[\rho \prod_{j=0}^{N-1} (a_{0,[j]} + a_{3,[j]}^{(n_j)} \sigma_{3,[j]})]$
probabilities $\uparrow \uparrow \uparrow$
 $\% \ \% \ \%$
measured: $\tilde{P} \rightarrow P$
 $\sum_{\vec{m}} M_{\vec{n}; \vec{m}} P_{\vec{m}}$
ideal \downarrow *ideal* \uparrow *already measured*
coefficients $M_{\vec{n}; \vec{m}}$
calculated $\sum_{\vec{m}} |P_{\vec{m}} - \tilde{P}_{\vec{m}}|^2$
 $P_{\vec{m}} = M^{-1} \tilde{P}_{\vec{m}}$
 \downarrow *constran:* $P_{\vec{m}} \geq 0$

- Can invert ideal measurement distribution by inversion (under constraint P be non-negative) [mentioned in an earlier lecture]

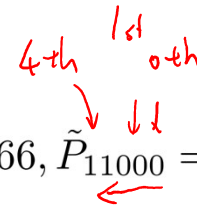
Simple application of mitigation

- Task: create a Bell state and measure in 0/1 basis



$$\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \otimes |000\rangle$$

- Results---the largest two components: $\tilde{P}_{00000} = 0.466, \tilde{P}_{11000} = 0.422$
 Smaller ones: $\tilde{P}_{00001} = 0.013, \tilde{P}_{01000} = 0.042, \tilde{P}_{10000} = 0.032, \tilde{P}_{11001} = 0.011$



- Invert to find P with constraint $P \geq 0$: $\tilde{P}_{\vec{n}} = \sum_{\vec{m}} M_{\vec{n};\vec{m}} P_{\vec{m}}$

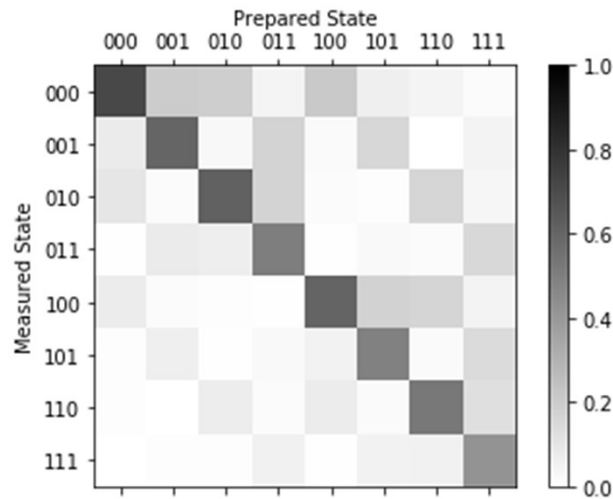
✓ The 'corrected' P has two dominant components, $P_{00000} = 0.493$ and $P_{11000} = 0.507$, with other components very small

- Can invert ideal measurement distribution by inversion (under constraint P be non-negative)

IBM has similar implementation

[qiskit-ignis/qiskit/advanced/ignis/ 4_measurement_error_mitigation.ipynb](https://github.com/qiskit-ignis/qiskit-ignis/blob/master/tutorials/4_measurement_error_mitigation.ipynb)

- ❑ Prepare and measure in computational basis



- Since one is concerned with only measurement distribution
- Obtain matrix M , then apply its inverse to mitigate errors

Other “tomographic” tools

- ❑ Disadvantage of current tomographic tools is the number of basis in the measurement scales exponentially with no. of qubits

- ❑ Randomized benchmarking

 - ➔ Quantify average error rates for gates

- ❑ Quantum volume

 - ➔ Quantify overall quality of quantum circuits

 - May return to these later (if time permits) and will do Qiskit demo