

Very long baselines with a superbeam

Very long baselines with a superbeam

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Neutrinos and New Physics

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FOR BNL Neutrino Working Group.

Wide Band Conventional Beam from BNL to the
Homestake Laboratory.

Summary of our study

- Baseline of > 2000 km with wide band conventional beams are the next step in accelerator neutrino physics.
- Extraordinary, large physical effects will be seen in such an experiment.
- Very good sensitivity to neutrino properties.
 - $< 1\%$ resolution on Δm_{32}^2
 - $< 1\%$ resolution on $\sin^2 2\theta_{23}$
 - Sensitivity to $\sin^2 2\theta_{13} > 0.005$ over a wide range of Δm_{32}^2
 - Sensitivity to CP violation.
 - Sign of Δm_{32}^2 over a wide range of parameters.
 - Measurement of Δm_{12}^2 in LMA region.
- Requires new thinking on how to build a beam and a detector. But experiment is technically feasible.

Comments

- Important ideas here are:

Long baseline to achieve large effects

Low energy wide band beam to get spectra

Beam is wide band, but low energy to make low backgrounds to ν_e appearance signature.

- Important difference between quark-matrix and neutrino-matrix

Neutrino oscillation effects are exactly calculable for any given set of parameters.
(including matter)

For quarks we often need complex tools such as CHPT and Lattice to connect CKM-matrix to physical phenomena.

- It makes sense to make a neutrino oscillation experiment with large effects even if they are sensitive to multiple parameters.

Neutrino Physics: the simple stuff

Assume a 2×2 neutrino mixing matrix.

$$\begin{pmatrix} \nu_a \\ \nu_b \end{pmatrix} = \begin{pmatrix} \cos(\theta) & \sin(\theta) \\ -\sin(\theta) & \cos(\theta) \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix} \quad (1)$$

$$\begin{aligned} \nu_a(t) &= \cos(\theta)\nu_1(t) + \sin(\theta)\nu_2(t) \\ P(\nu_a \rightarrow \nu_b) &= | \langle \nu_b | \nu_a(t) \rangle |^2 \\ &= \sin^2(\theta) \cos^2(\theta) |e^{-iE_2 t} - e^{-iE_1 t}|^2 \end{aligned} \quad (2)$$

Sufficient to understand most of the physics:

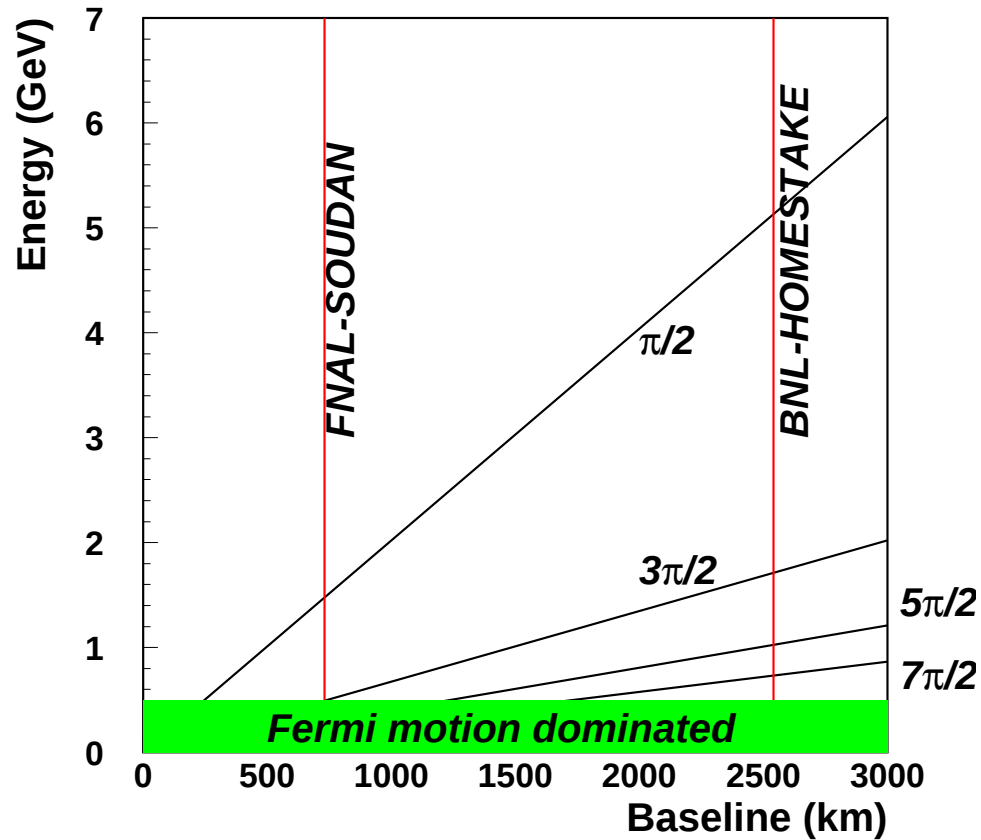
$$P(\nu_a \rightarrow \nu_b) = \sin^2 2\theta \sin^2 \frac{1.27(\Delta m^2/eV^2)(L/km)}{(E/GeV)}$$

$$P(\nu_a \rightarrow \nu_a) = 1 - \sin^2 2\theta \sin^2 \frac{1.27(\Delta m^2/eV^2)(L/km)}{(E/GeV)}$$

Oscillation nodes at $\pi/2, 3\pi/2, 5\pi/2, \dots$ ($\pi/2$):

$$\Delta m^2 = 0.003 eV^2, E = 1 GeV, L = 412 km .$$

Oscillation Nodes for $\Delta m^2 = 0.0025 \text{ eV}^2$



- Large effects: Multiple oscillation nodes.
- Fermi motion limits resolution at low energies: wide band beam ($0.5 \rightarrow 8 \text{ GeV}$).
- $\Delta m^2 \approx 0.0025 \text{ eV}^2$: Baseline $> 2000 \text{ km}$.

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Neutrino Physics: the difficult stuff

Bill Marciano, hep-ph/0108181

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} c_1 c_3 & s_1 c_3 & s_3 e^{-i\delta} \\ -s_1 c_2 - c_1 s_2 s_3 e^{i\delta} & c_1 c_2 - s_1 s_2 s_3 e^{i\delta} & s_2 c_3 \\ s_1 s_2 - c_1 c_2 s_3 e^{i\delta} & -c_1 s_2 - s_1 c_2 s_3 e^{i\delta} & c_2 c_3 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} \quad (3)$$

$$\begin{aligned} P(\nu_\mu \rightarrow \nu_e) = & 4(s_2^2 s_3^2 c_3^2 + J_{CP} \sin \Delta_{21}) \sin^2 \frac{\Delta_{31}}{2} \\ & + 2(s_1 s_2 s_3 c_1 c_2 c_3^2 \cos \delta - s_1^2 s_2^2 s_3^2 c_3^2) \sin \Delta_{31} \sin \Delta_{21} \\ & + 4(s_1^2 c_1^2 c_2^2 c_3^2 + s_1^4 s_2^2 s_3^2 c_3^2 - 2s_1^3 s_2 s_3 c_1 c_2 c_3^2 \cos \delta \\ & - J_{CP} \sin \Delta_{31}) \sin^2 \frac{\Delta_{21}}{2} \\ & + 8(s_1 s_2 s_3 c_1 c_2 c_3^2 \cos \delta - s_1^2 s_2^2 s_3^2 c_3^2) \sin^2 \frac{\Delta_{31}}{2} \sin^2 \frac{\Delta_{21}}{2} \end{aligned} \quad (4)$$

No matter effects in above formula

$$\begin{aligned}\Delta_{31} &\equiv \Delta m_{31}^2 L/2E_\nu \\ \Delta_{21} &\equiv \Delta m_{21}^2 L/2E_\nu\end{aligned}$$

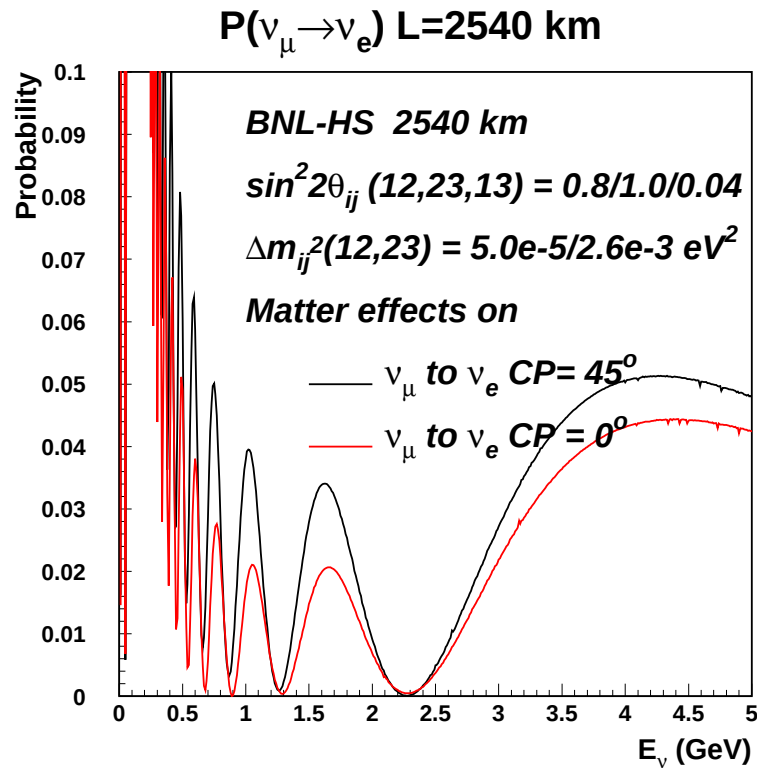
$$J_{CP} \equiv s_1 s_2 s_3 c_1 c_2 c_3^2 \sin \delta \quad (5)$$

$$A \equiv \frac{P(\nu_\mu \rightarrow \nu_e) - P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)}{P(\nu_\mu \rightarrow \nu_e) + P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)} \quad (6)$$

To leading order in Δ_{21} (assumed to be small), one finds

$$P(\nu_\mu \rightarrow \nu_e) \simeq 4s_2^2 s_3^2 c_3^2 \sin^2 \frac{\Delta_{31}}{2} + \mathcal{O}(\Delta_{21}) \quad (12a)$$

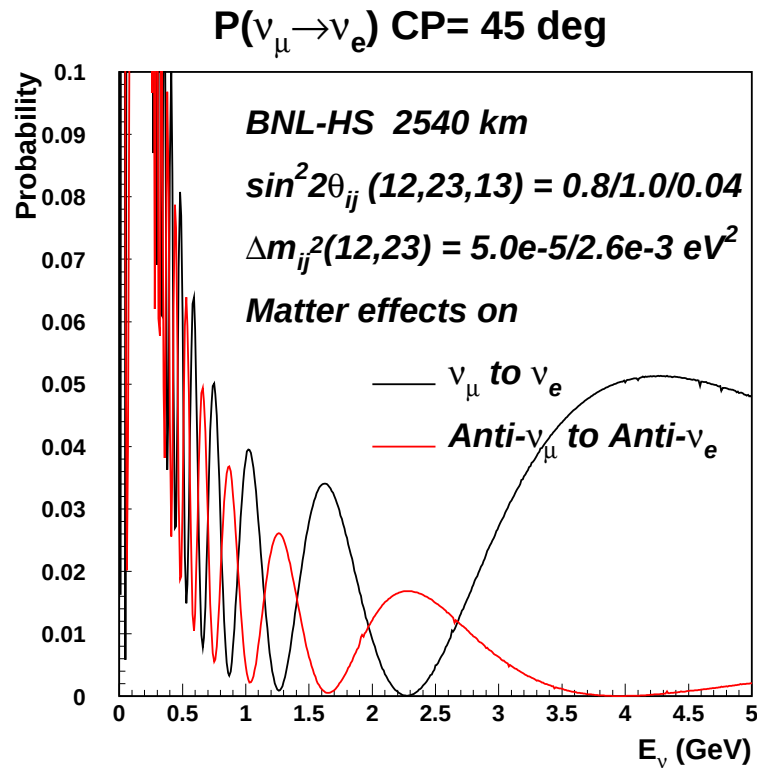
$$A \simeq \frac{J_{CP} \sin \Delta_{21}}{s_2^2 s_3^2 c_3^2} \simeq \frac{2s_1 c_1 c_2 \sin \delta}{s_2 s_3} \left(\frac{\Delta m_{21}^2}{\Delta m_{31}^2} \right) \frac{\Delta m_{31}^2 L}{4E_\nu} + \mathcal{O}(\Delta_{21}^2) \quad (12b)$$



General Features

- 0.5 – 1 GeV: Δm_{12}^2 (LMA) region.
- 1 – 3 GeV: CP large effects region
- > 3 GeV: Matter enhanced (ν_μ), suppressed ($\bar{\nu}_\mu$). ($\Delta m_{32}^2 > 0$) Region.

I. Mocioiu and R. Shrock, Phys. Rev. D62, 053017 (2000), JHEP 0111, 050 (2001)



Compare Neutrino to Antineutrino.

- 0.5 – 1 GeV: Δm_{12}^2 (LMA) region.
- 1 – 3 GeV: CP region
- > 3 GeV: Matter enhanced (ν_μ), suppressed ($\bar{\nu}_\mu$). ($\Delta m_{32}^2 > 0$) Region.

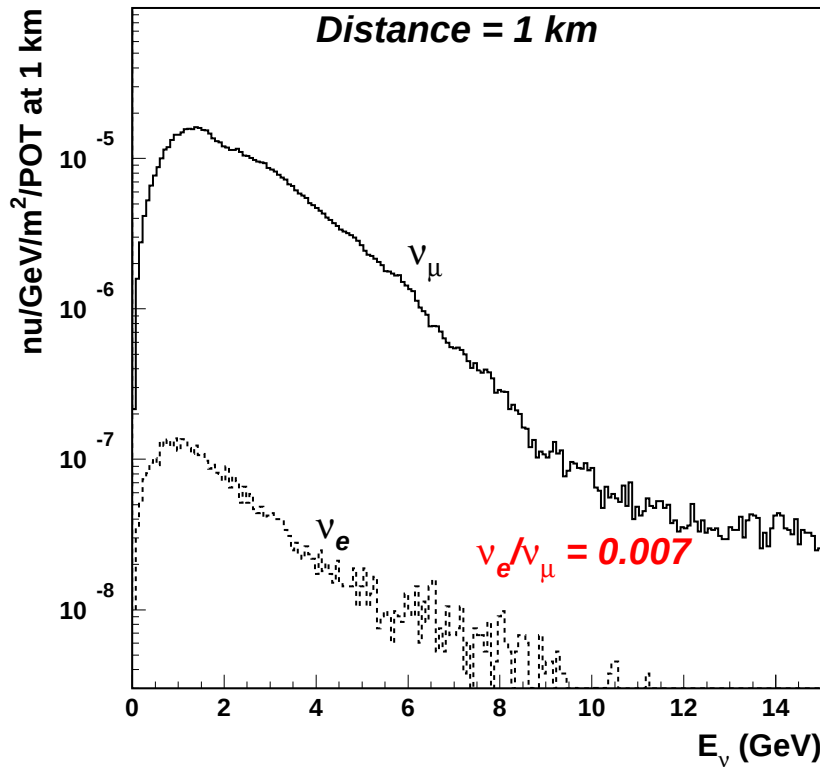
Very long baselines with a superbeam

4 GOALS OF NEUTRINO OSCILLATION PHYSICS

- Precise determination of Δm_{32}^2 and $\sin^2 2\theta_{23}$ and definitive observation of oscillatory behavior.
- Detection of $\nu_\mu \rightarrow \nu_e$ in the appearance mode. If $\Delta m_{\nu_\mu \rightarrow \nu_e}^2 = \Delta m_{32}^2$ then $|U_{e3}|^2 (= \sin^2 \theta_{13})$ is non-zero.
- Detection of the matter enhancement effect in $\nu_\mu \rightarrow \nu_e$. Sign of Δm_{32}^2 ; i.e. which neutrino is heavier.
- Detection of CP violation in neutrino physics. Phase of $|U_{e3}|$ is CP violating and causes asymmetry in the rates $\nu_\mu \rightarrow \nu_e$ versus $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$.

It will be good to do it all in same experiment with only neutrino beam (no antineutrino).

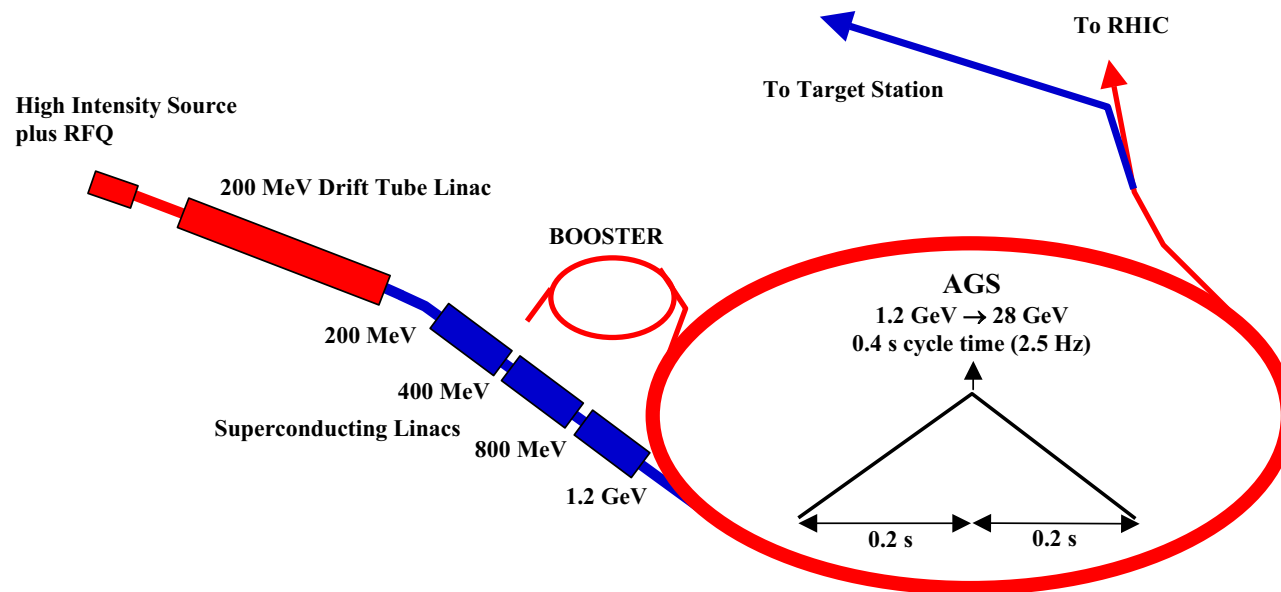
BNL Wide Band. Proton Energy = 28 GeV



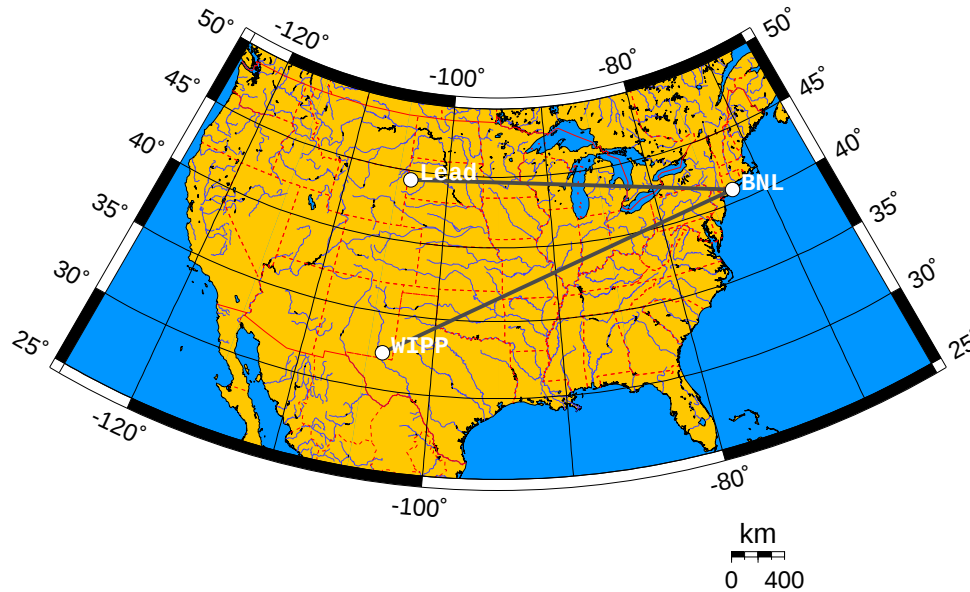
- New design spans 0.5-6 GeV
- Low ν_e background 0.7%
 0.0073 ± 0.0014 (E734 1986).
- Low background from high energies (NC and ν_τ for ν_e)
- 200 m decay tunnel
- Graphite target embedded in horn
- Target cooling achievable for 1 MW

The Accelerator

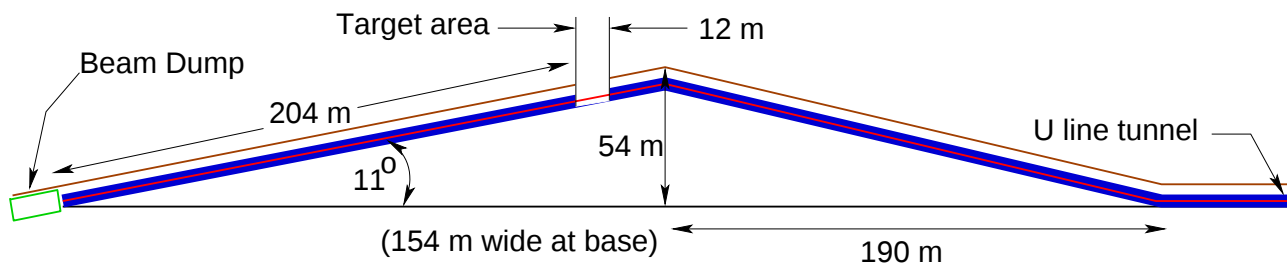
- Conceptually simple upgrade. No magic. Cost $\sim \$100M$.
- Run 28 GeV AGS at 2.5 Hz to get 1 MW.
- Need faster proton source: Super Conducting LINAC at 1.2 GeV
- Current: $7 \times 10^{13} ppp$ at 0.5 Hz $\Rightarrow$ LINAC: $10^{14} ppp$ at 2.5 Hz.



Beam on the Hill



- BNL-Lead 2540km
BNL-Wipp: 2880km
- Avoids water table.
- Hills are inexpensive:
highway ramps.
- Total cost \$35 M for
200 m tunnel.



Very long baselines with a superbeam

Event Rates with Neutrinos

Assume 1 MW, 500 kT Fiducial, 5×10^7 sec running. (1.22×10^{22} Protons at 28 GeV.)

Assume Water Cerenkov detector (with $\sim 10\%$ PMT coverage)

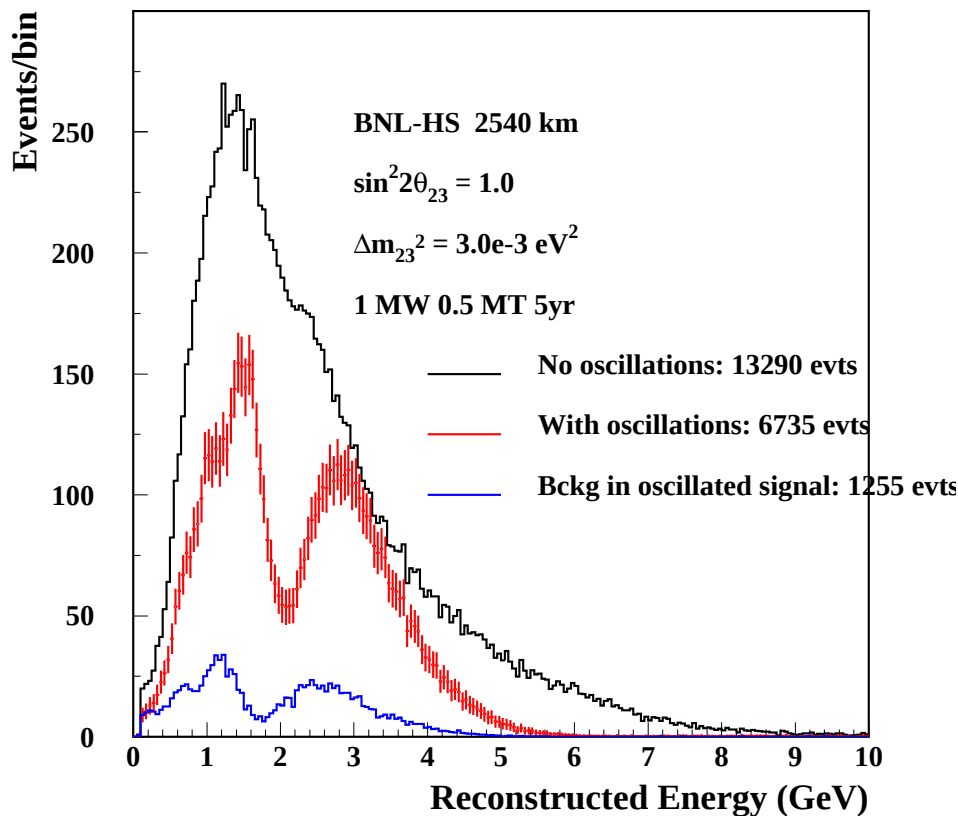
CC $\nu_\mu + N \rightarrow \mu^- + X$	51800
NC $\nu_\mu + N \rightarrow \nu_\mu + X$	16908
CC $\nu_e + N \rightarrow e^- + X$	380
QE $\nu_\mu + n \rightarrow \mu^- + p$	11767
QE $\nu_e + n \rightarrow e^- + p$	84
CC $\nu_\mu + N \rightarrow \mu^- + \pi^+ + N$	14574
NC $\nu_\mu + N \rightarrow \nu_\mu + N + \pi^0$	3178
NC $\nu_\mu + O^{16} \rightarrow \nu_\mu + O^{16} + \pi^0$	574
CC $\nu_\tau + N \rightarrow \tau^- + X$ (if all $\nu_\mu \rightarrow \nu_\tau$)	319

Backgrounds to clean (QE) events SMALL

NC dominated by elastic and single π .

Low τ production.

ν_μ DISAPPEARANCE

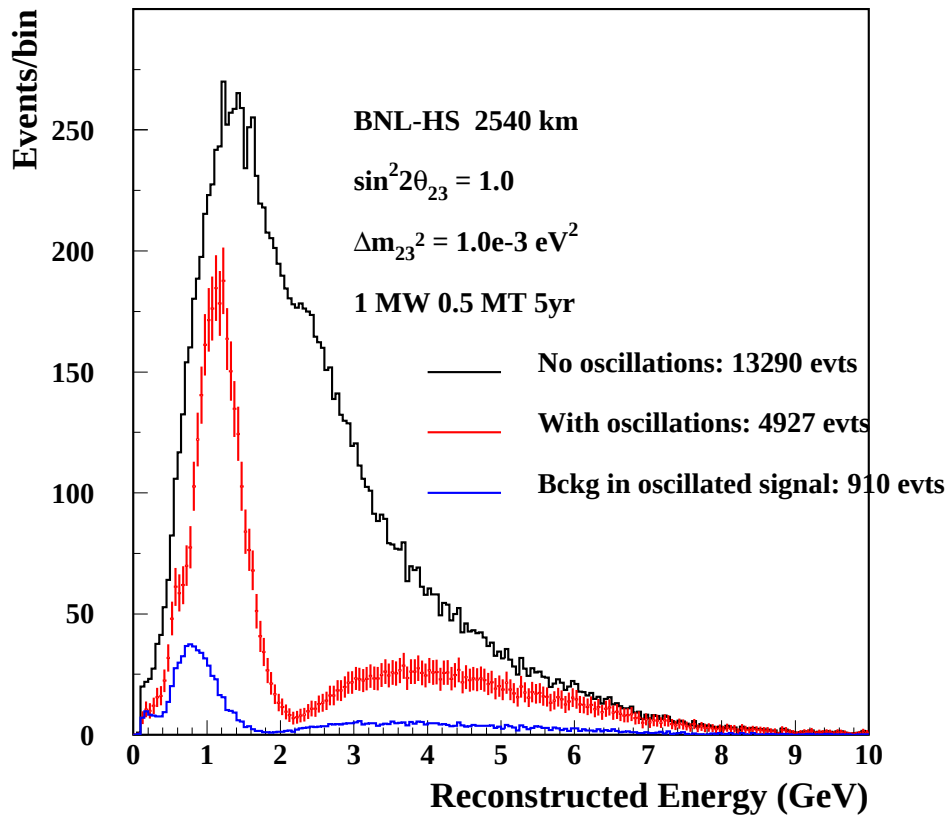


Node pattern provides high Δm_{23}^2 resolution.
 Energy calibration is very important.

Flux normalization not important for
 measurement of $\sin^2 2\theta_{23}$

Background shape can be measured independently
 Minimum systematics in ν_μ and $\bar{\nu}_\mu$ comparison

ν_μ DISAPPEARANCE

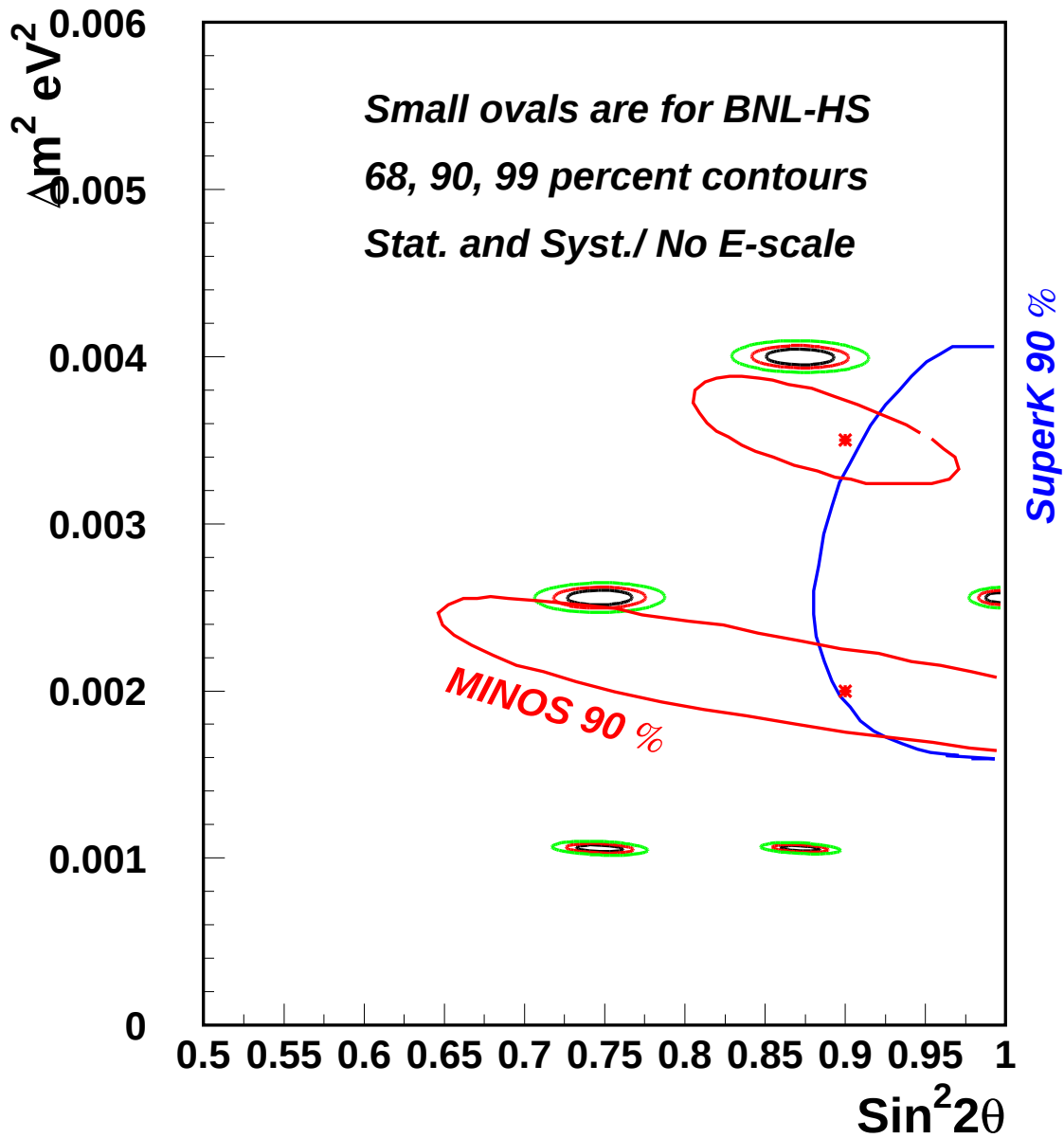


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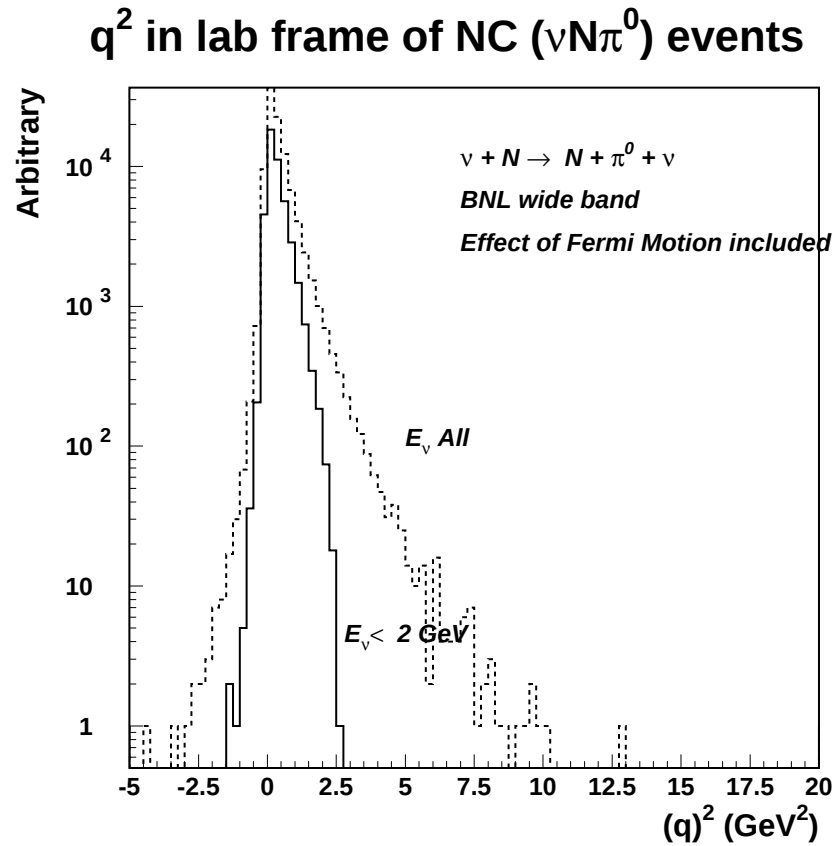
Test points for ν_μ disapp



Measurement of Δm_{23}^2

- Little dependence on systematic errors on resolution, backgrounds, energy linearity, or normalization.
- Ultimate resolution on Δm^2 depends on energy calibration. For perfect energy calibration $\pm 0.7\%$ possible.
- Energy calibration at $< 1\%$ in 1-5 GeV region needed.
- Can exclude $\sin^2 2\theta_{23} < 0.99$ at 90% C.L. Could be better with accurate background subtraction.
- No need of near detector for this measurement. Even a 10% systematic error on normalization does not bother measurement.

NC π^0 background for $\nu_\mu \rightarrow \nu_e$



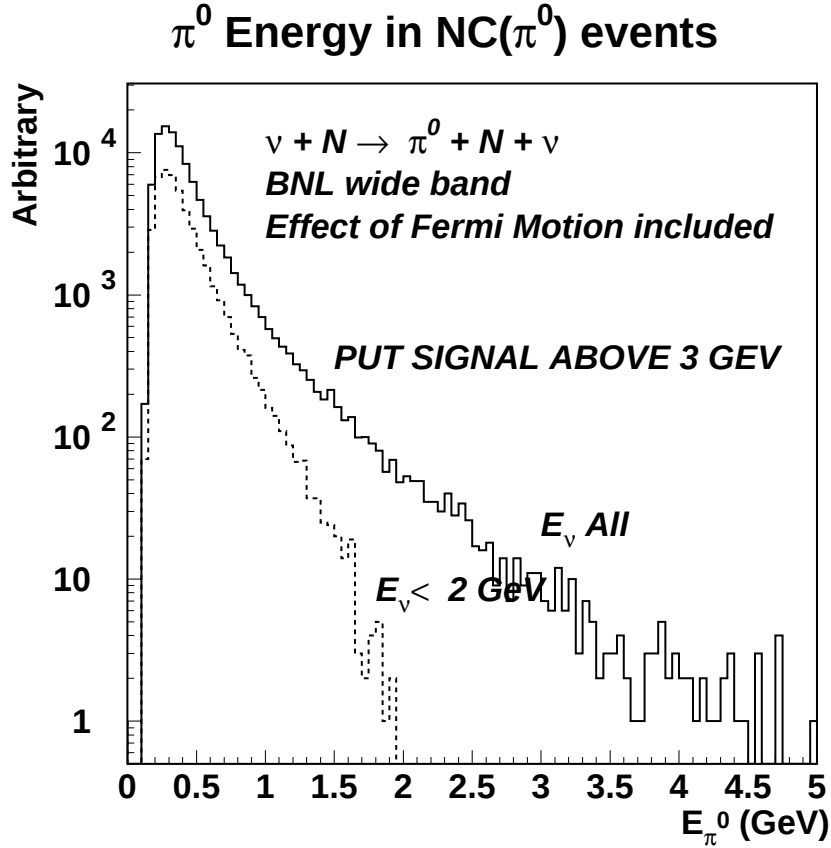
$$q^2 = (p'_N + p'_\pi) - p_N.$$

$p_N = (0, 0, 0, m_N)$ At rest nucleon.

General feature of all neutral current processes:

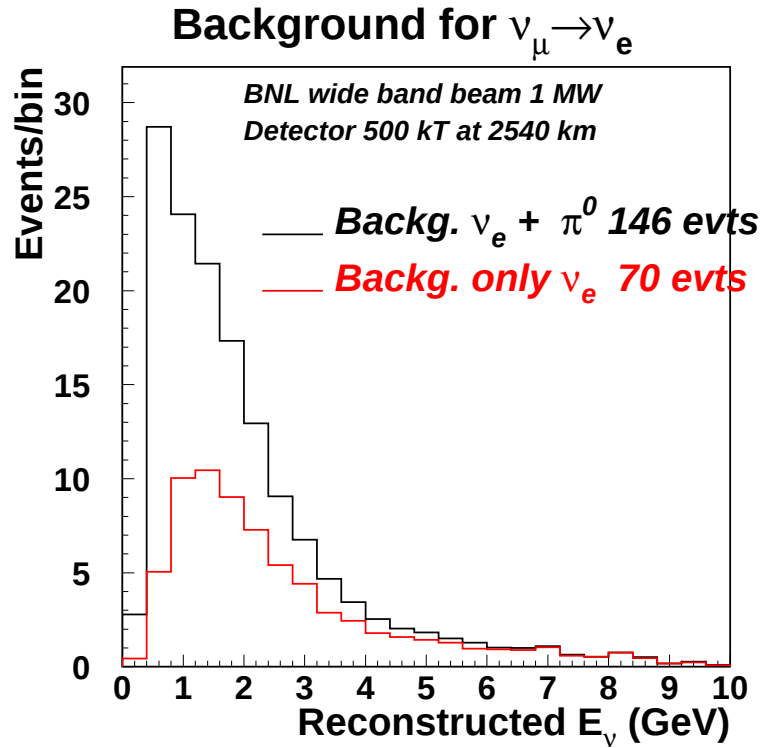
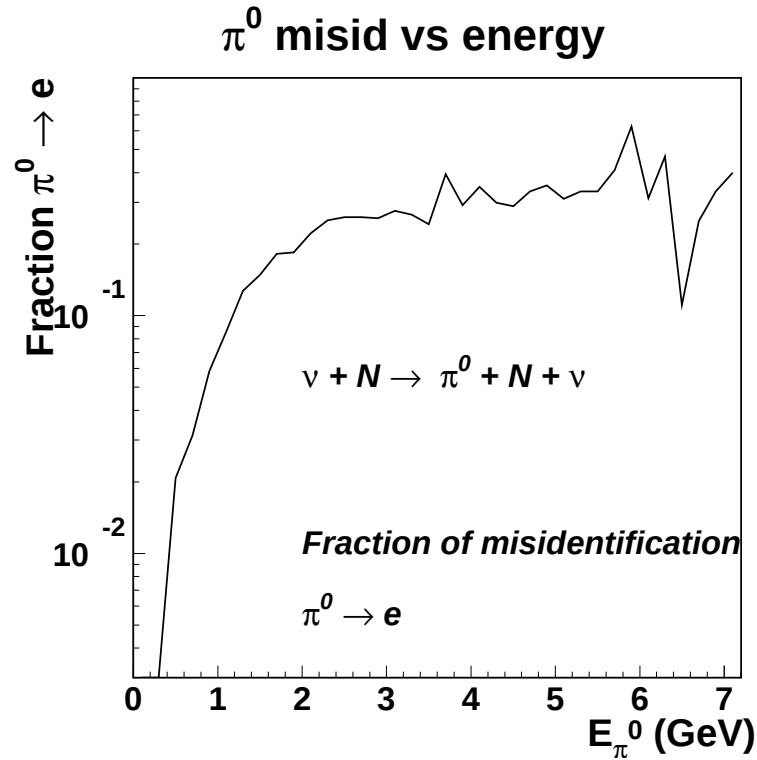
Low q^2 or low hadronic energy in final state
 independent of neutrino energy.

NC π^0 background for $\nu_\mu \rightarrow \nu_e$



- The NC energy distribution is independent of ν -energy except the kinematic limit.
- In $\nu_\mu N \rightarrow \nu_\mu N \pi^0$ events all energy ν produce peak at the same energy except the tail.
- For a very long baselines and wide band beam ν_e signal will be above 3 GeV with little π^0 background.

$\nu_\mu \rightarrow \nu_e$ All background

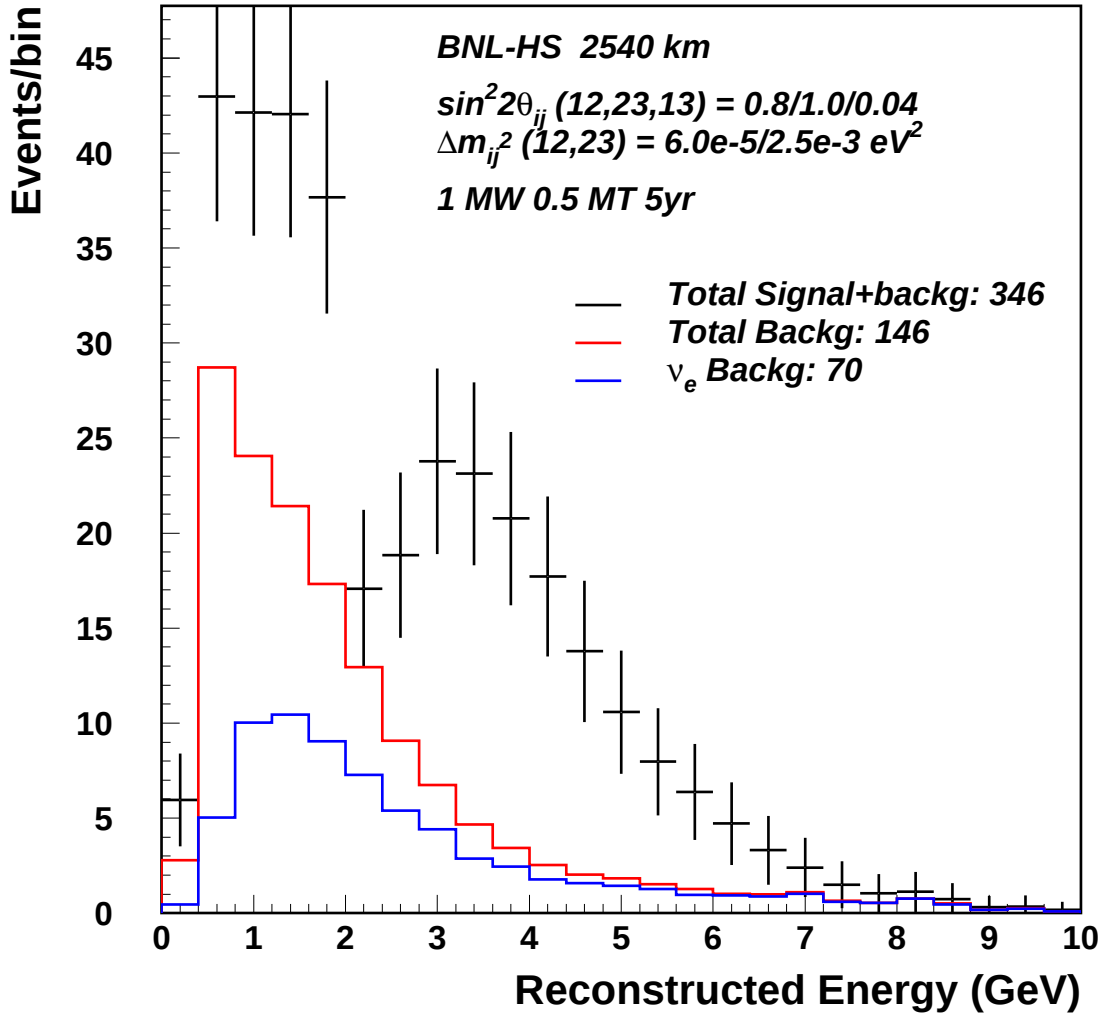


- Background includes $\nu N \pi^0$ and Coherent $\nu O^{16} \pi^0$.
- Efficiency for signal is $\sim 80\%$
- For $E_\nu < 2\text{GeV}$ $N_{\pi^0} : N_{\nu_e} :: 59 : 35$
- For $E_\nu > 2\text{GeV}$ $N_{\pi^0} : N_{\nu_e} :: 17 : 35$

Very long baselines with a superbeam

Measurement of $\sin^2 2\theta_{13}$

ν_e APPEARANCE



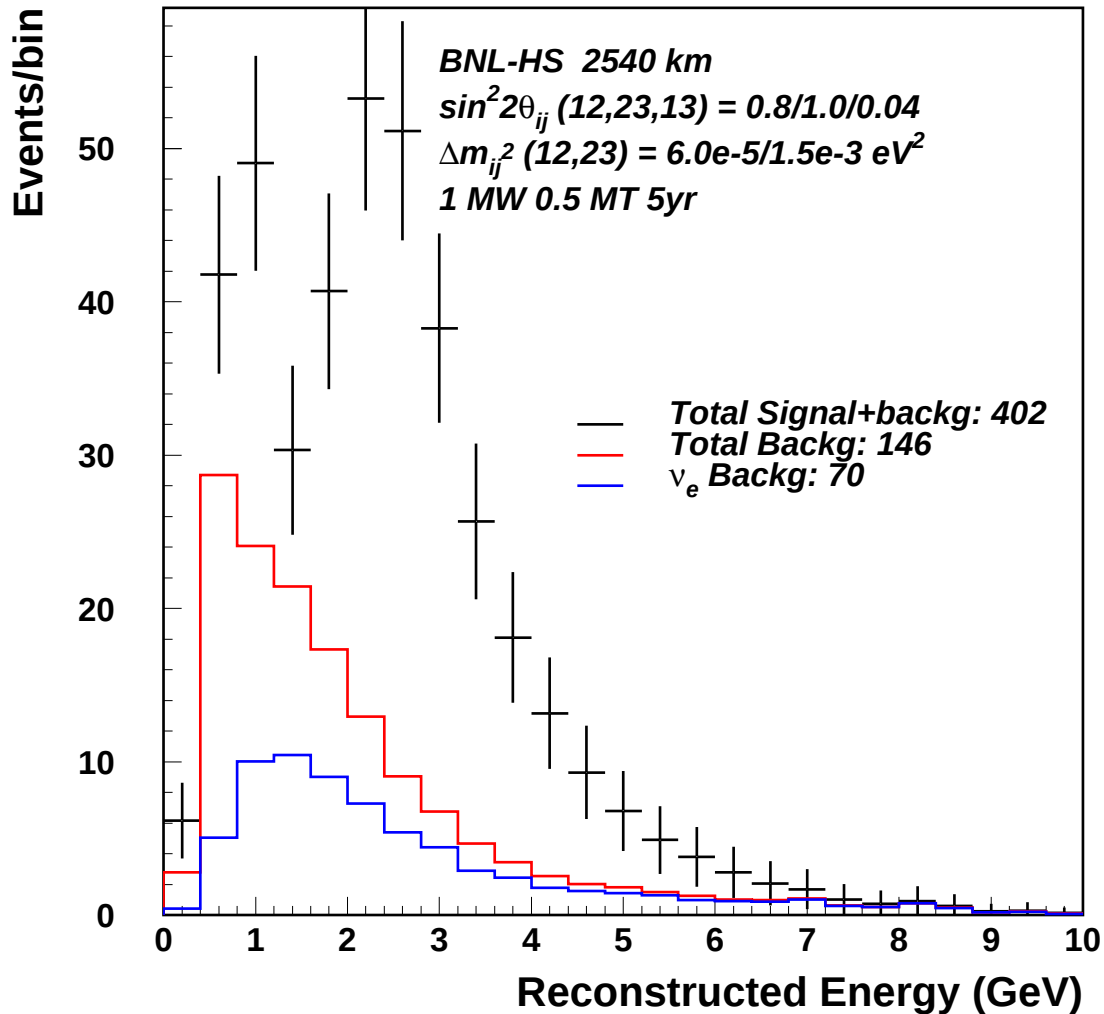
$$\Delta m_{23}^2 = 0.0025 \text{ eV}^2, \sin^2 2\theta_{13} = 0.04.$$

Assume normal mass hierarchy. $m_3 > m_2 > m_1$

Matter effects included.

Measurement of $\sin^2 2\theta_{13}$

ν_e APPEARANCE

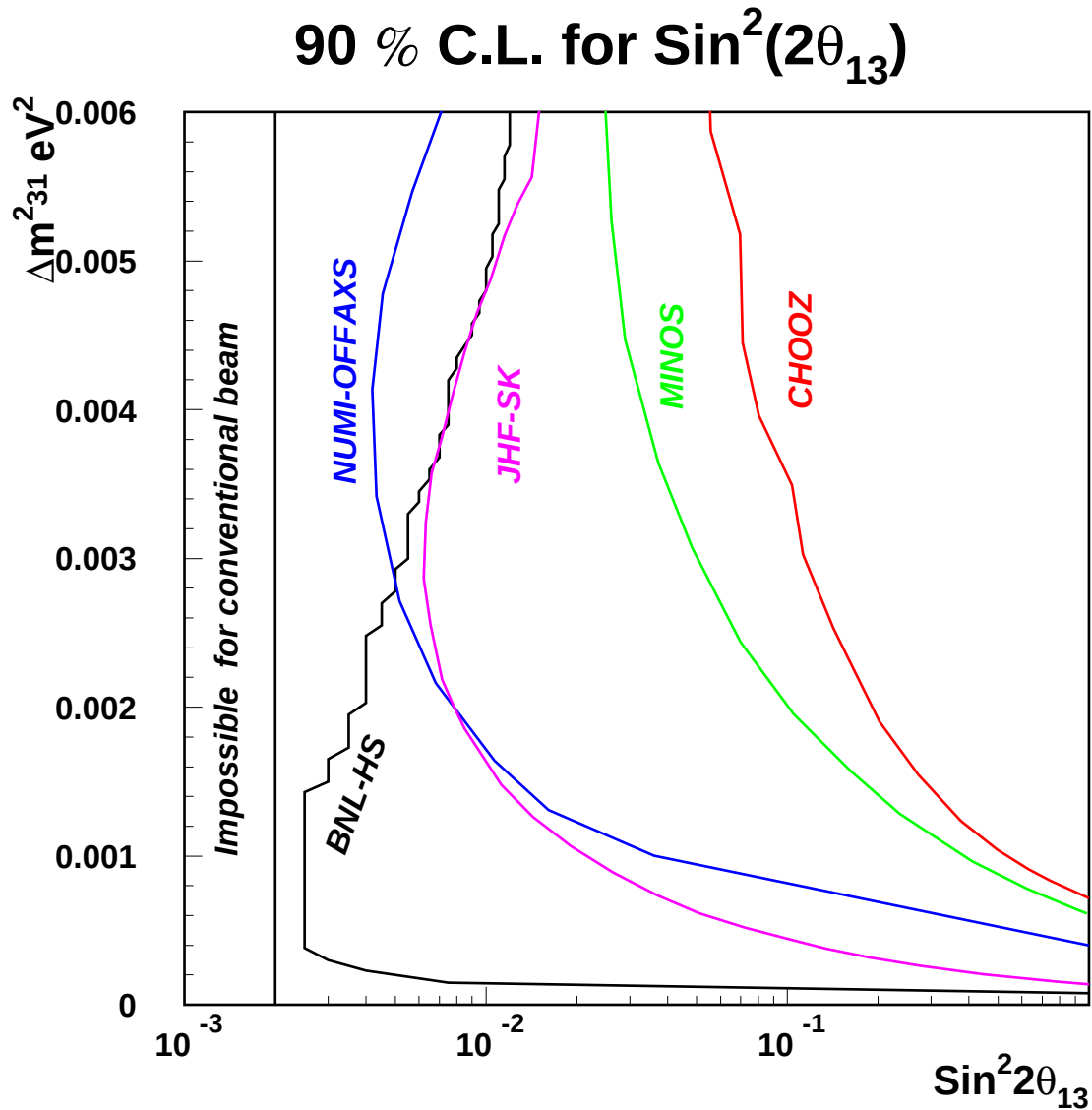


$$\Delta m_{23}^2 = 0.0015 \text{ eV}^2, \sin^2 2\theta_{13} = 0.04.$$

Assume normal mass hierarchy. $m_3 > m_2 > m_1$

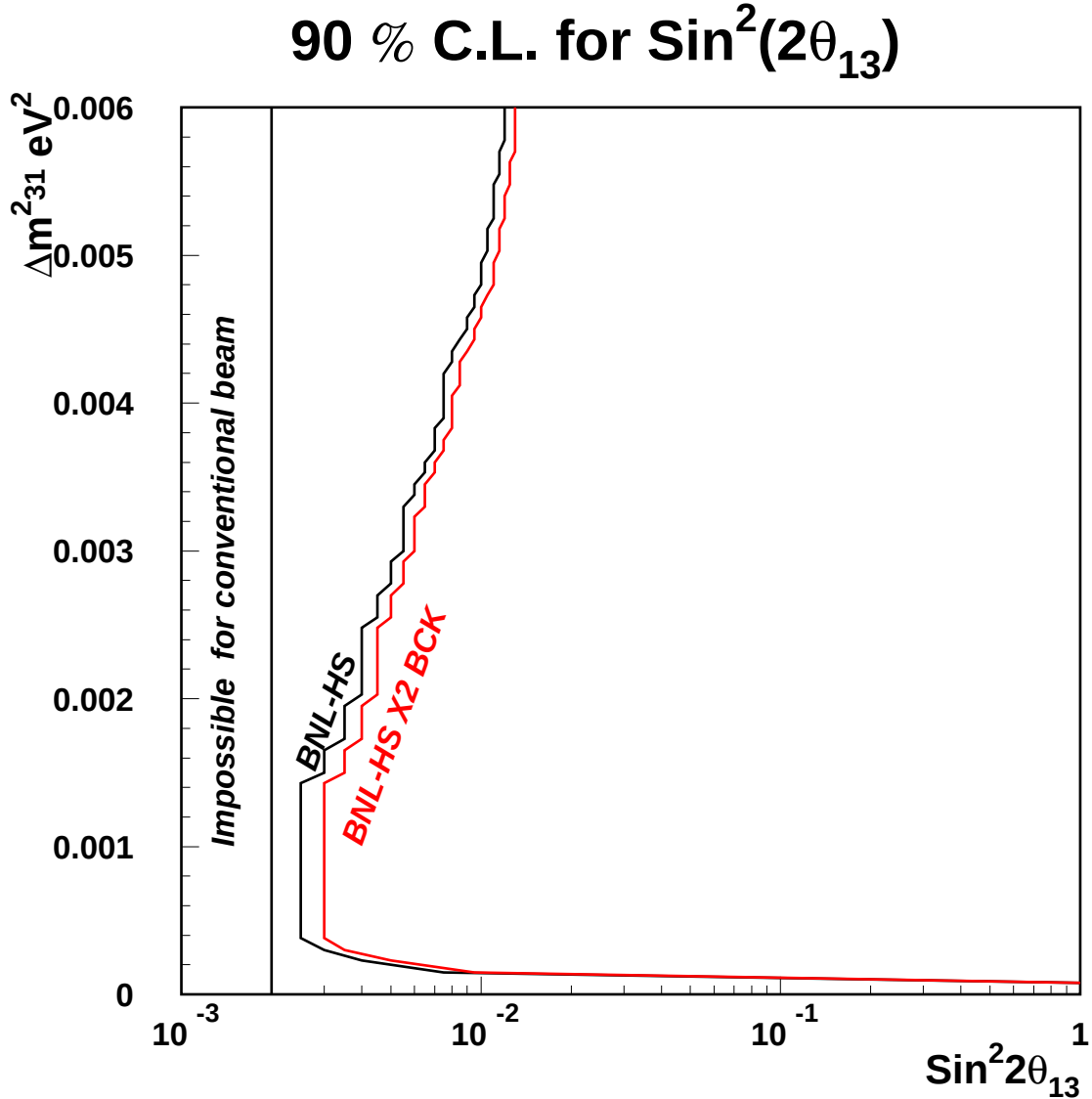
Matter effects included.

Measurement of $\sin^2 2\theta_{13}$ 90% C.L.



Distinctive signature with multiple oscillations
above 0.001 eV^2

Measurement of $\sin^2 2\theta_{13}$ 90% C.L. high Bckg.



Assume that the neutral current background is higher by factor of 2 over the entire spectrum.

Measurement of $\sin^2 2\theta_{13}$ 90% C.L.

BNL-HS(2540 km) good sensitivity to $\sin^2 2\theta_{13}$.

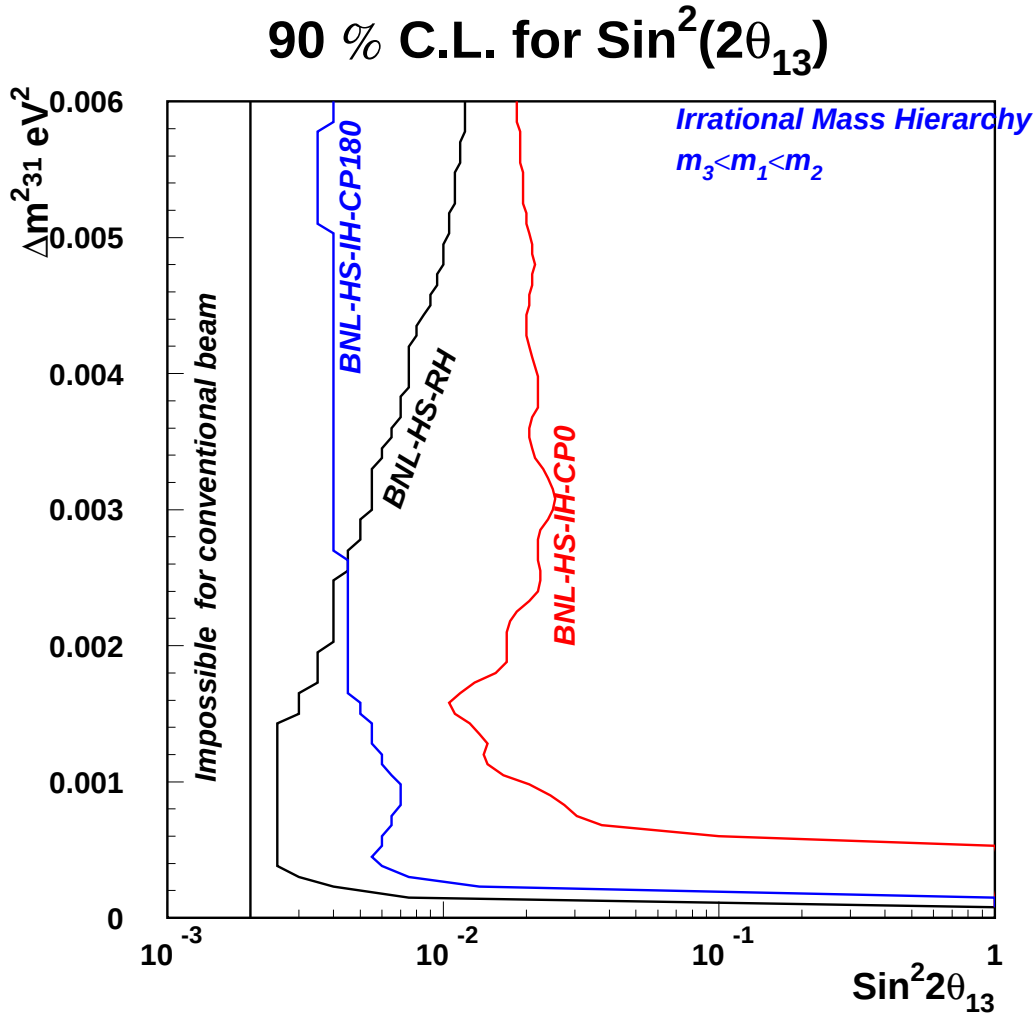
Improvement from 0.12 to 0.005 at 0.0025 eV^2 .

Signal very distinctive above 0.001 eV^2 .

Need harder beam to improve sensitivity above 0.004 eV^2 .

No experiment can go below $\sin^2 2\theta_{13} \approx 0.002$ with horn focussed beam due to systematic error on intrinsic ν_e background.

Mass Hierarchy



Regular Mass hierarchy: $m_3 > m_2 > m_1$ (RH)

Reversed Mass hierarchy: $m_1 > m_2 > m_3$ (RVH)

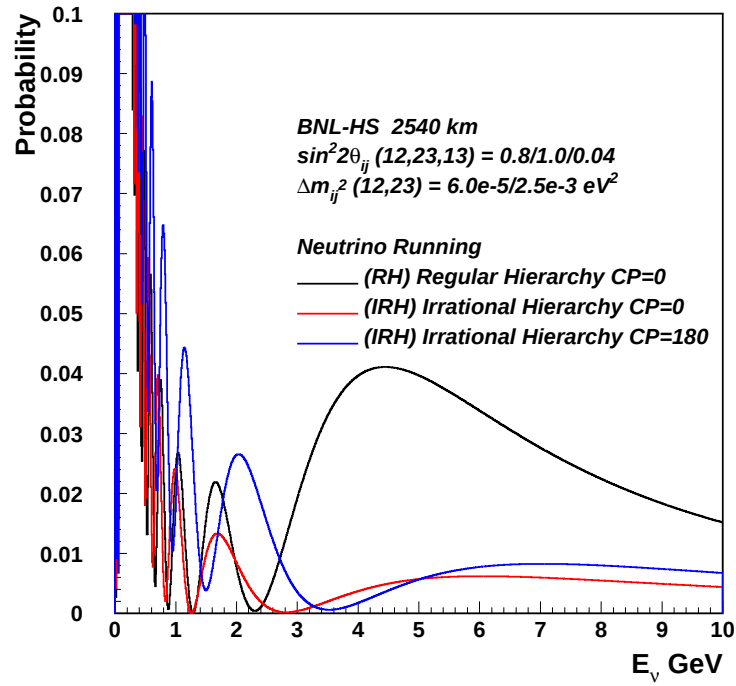
Irrational Mass hierarchy: $m_2 > m_1 > m_3$ (IH)

RVH is ruled out if Solar LMA is the correct solution.

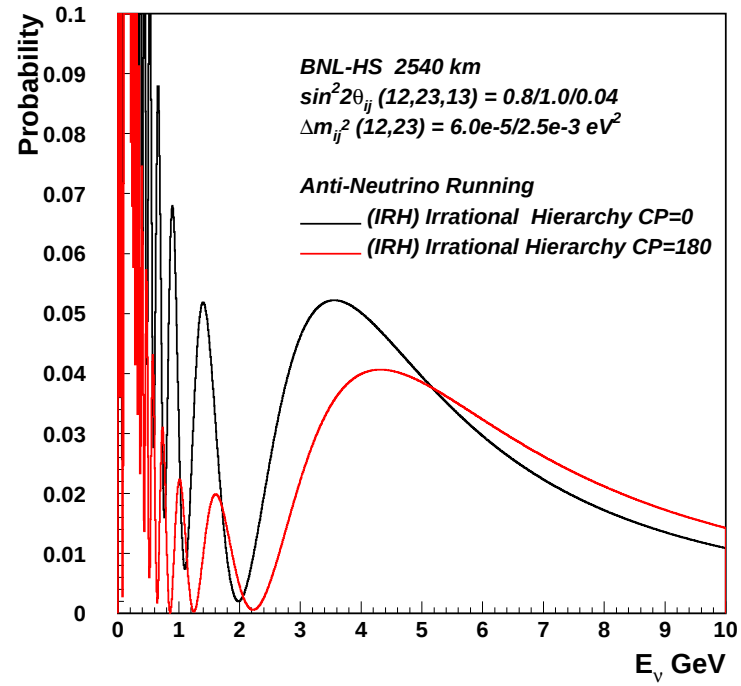
We would need to run Anti-neutrino beam to fully explore IH.

Mass Hierarchy Anti-neutrinos

$\nu_\mu \rightarrow \nu_e$ Oscillation

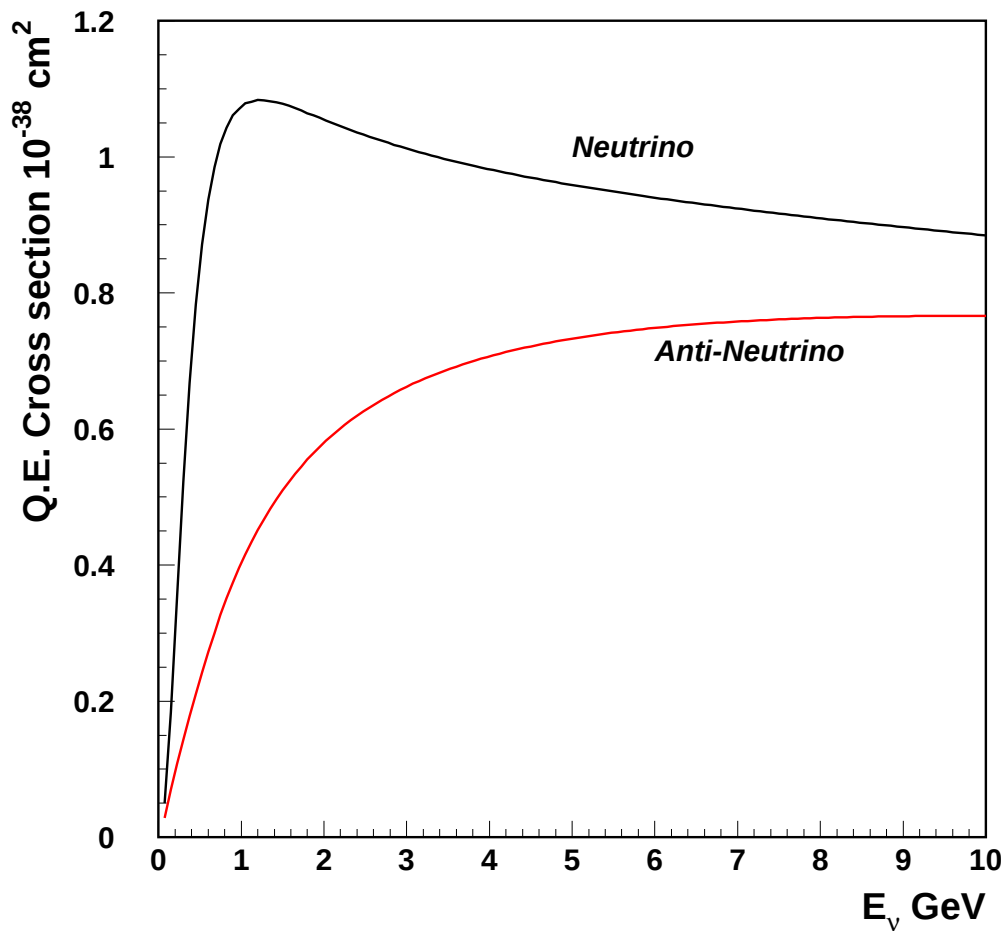


Anti- $\nu_\mu \rightarrow$ Anti- ν_e Oscillation



Very long baselines with a superbeam

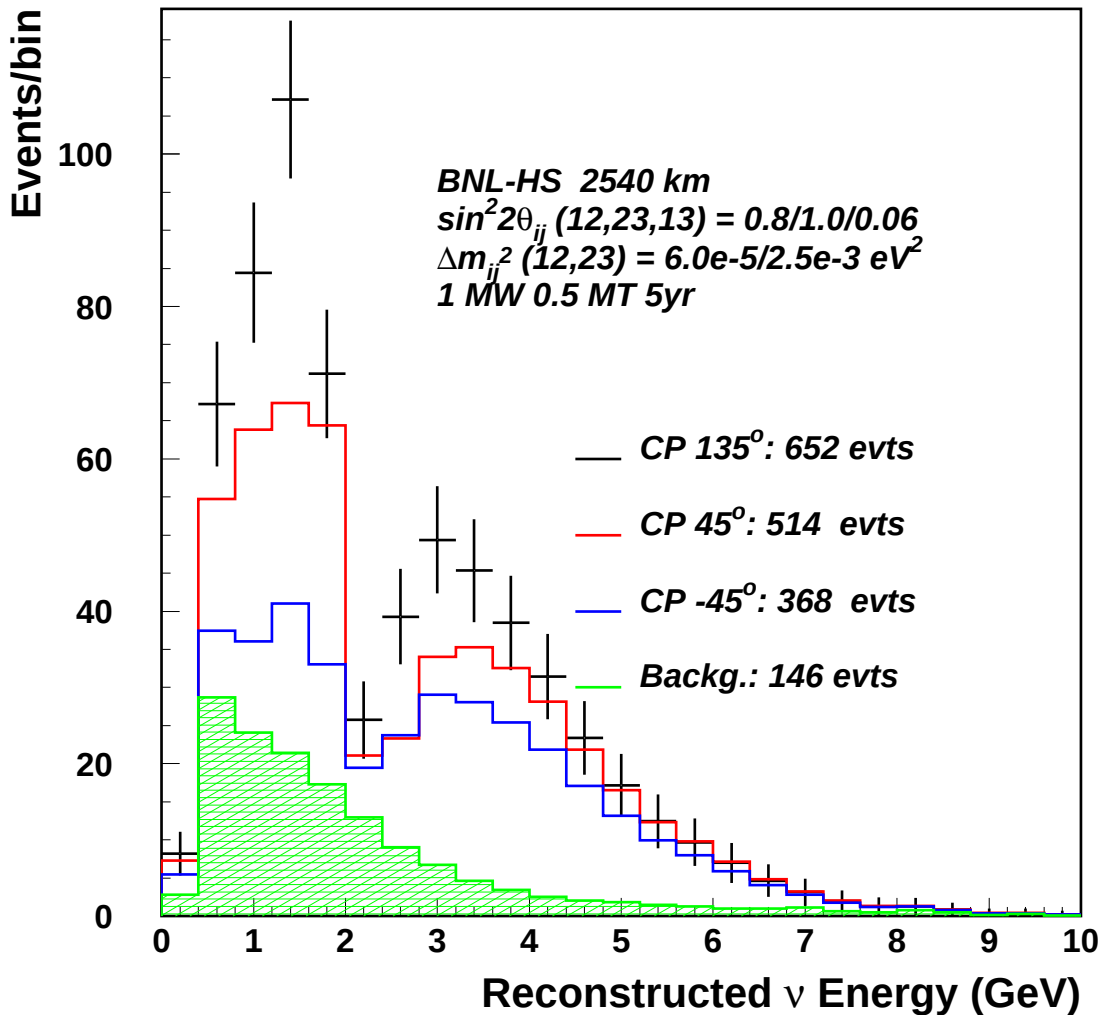
Quasielastic cross section



An experiment searching for signal at high energies may not need much more anti-neutrino running than neutrino running.

δ_{CP} Measurement. BNL-to-HS, 2540 km, 1 MW, 500kT, 5×10^7 sec

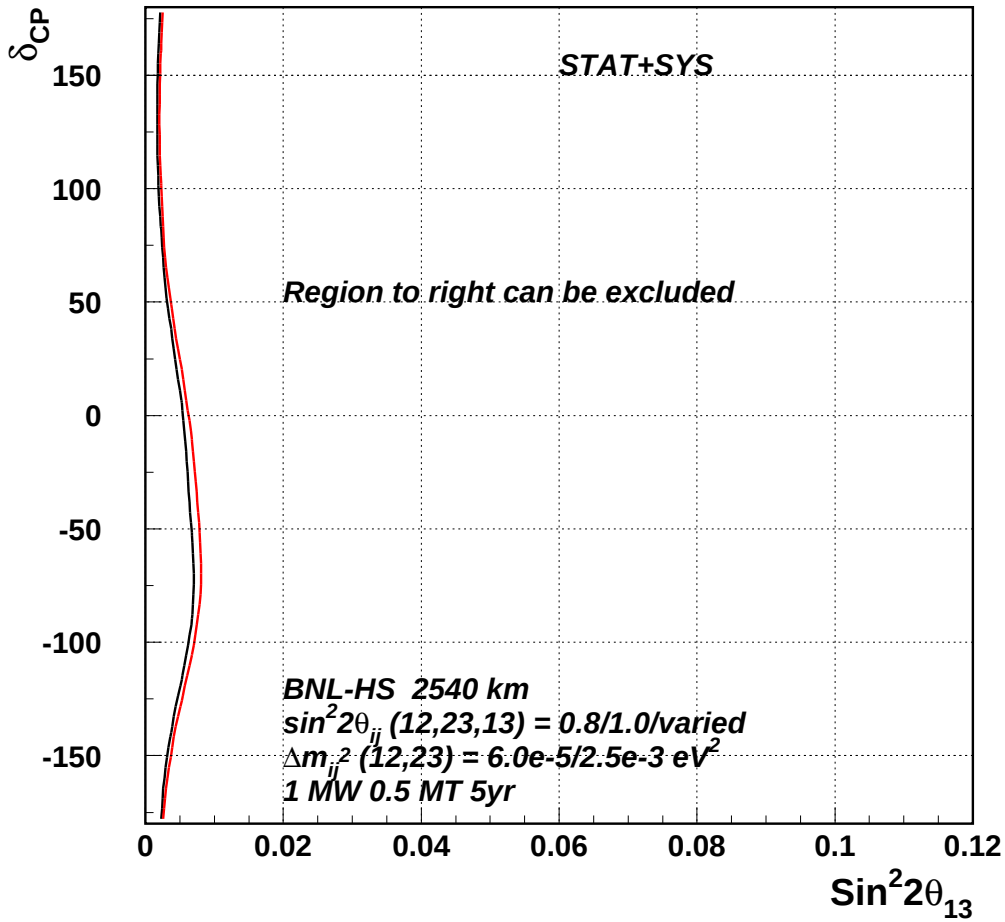
ν_e APPEARANCE



CP parameter can be determined from only neutrino data.
Good background subtraction can help.

Measurement of δ_{CP} ; Confidence Levels

90, 95 % C.L. for δ_{CP} vs $\sin^2 2\theta_{13}$



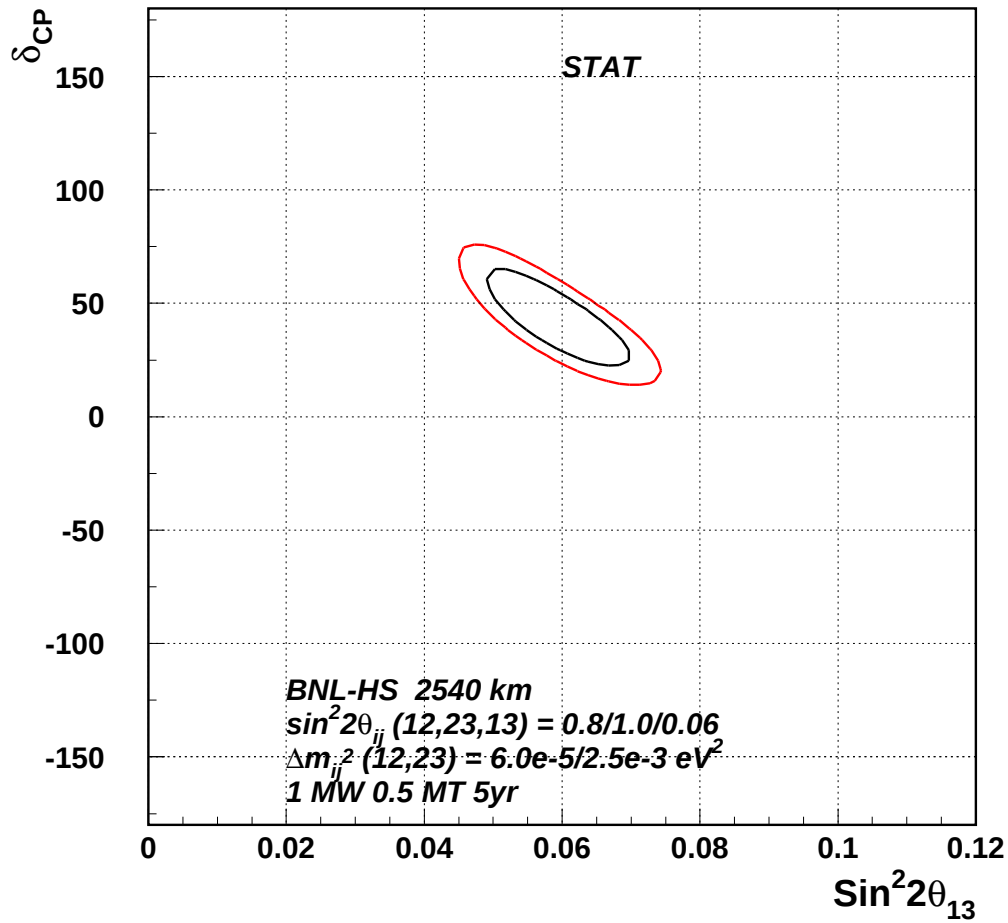
$$\Delta m_{21}^2 = 6 \times 10^{-5} \text{ eV}^2, \Delta m_{31}^2 = 2.5 \times 10^{-3} \text{ eV}^2$$

$$\sin^2 2\theta_{12} = 0.8, \sin^2 2\theta_{23} = 1.0$$

The region on the right hand side of curve can be excluded at 95% C.L.

Measurement of $\delta_{CP} = 45^\circ$

90 % C.L. for δ_{CP} vs $\sin^2 2\theta_{13}$



No Systematic error

$$\Delta m_{21}^2 = 6 \times 10^{-5} \text{ eV}^2, \Delta m_{31}^2 = 2.5 \times 10^{-3} \text{ eV}^2$$

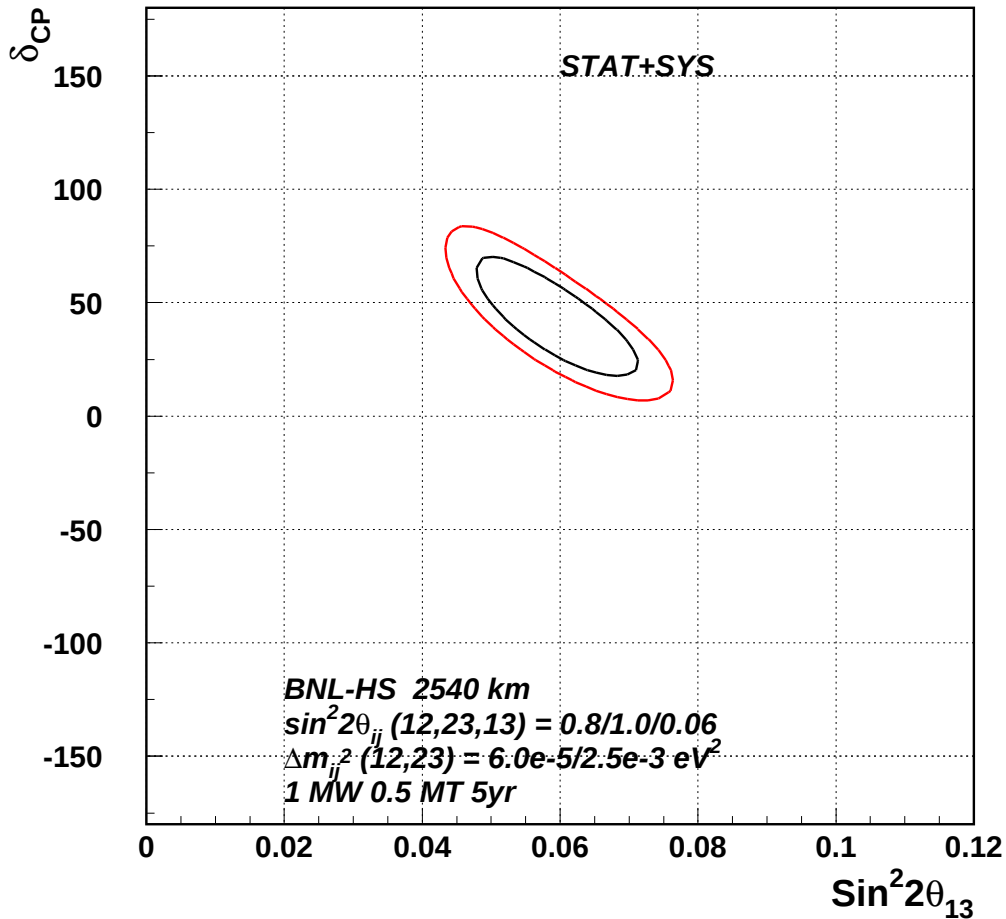
$$\sin^2 2\theta_{12} = 0.8, \sin^2 2\theta_{23} = 1.0$$

$$\delta_{CP} = 45^\circ, \sin^2 2\theta_{13} = 0.06$$

68%, and 90% C.L.

Measurement of $\delta_{CP} = 45^\circ$ No anti-neutrino running.

90 % C.L. for δ_{CP} vs $\sin^2 2\theta_{13}$



Systematic error of 10% on backg.

$$\Delta m_{21}^2 = 6 \times 10^{-5} \text{ eV}^2, \Delta m_{31}^2 = 2.5 \times 10^{-3} \text{ eV}^2$$

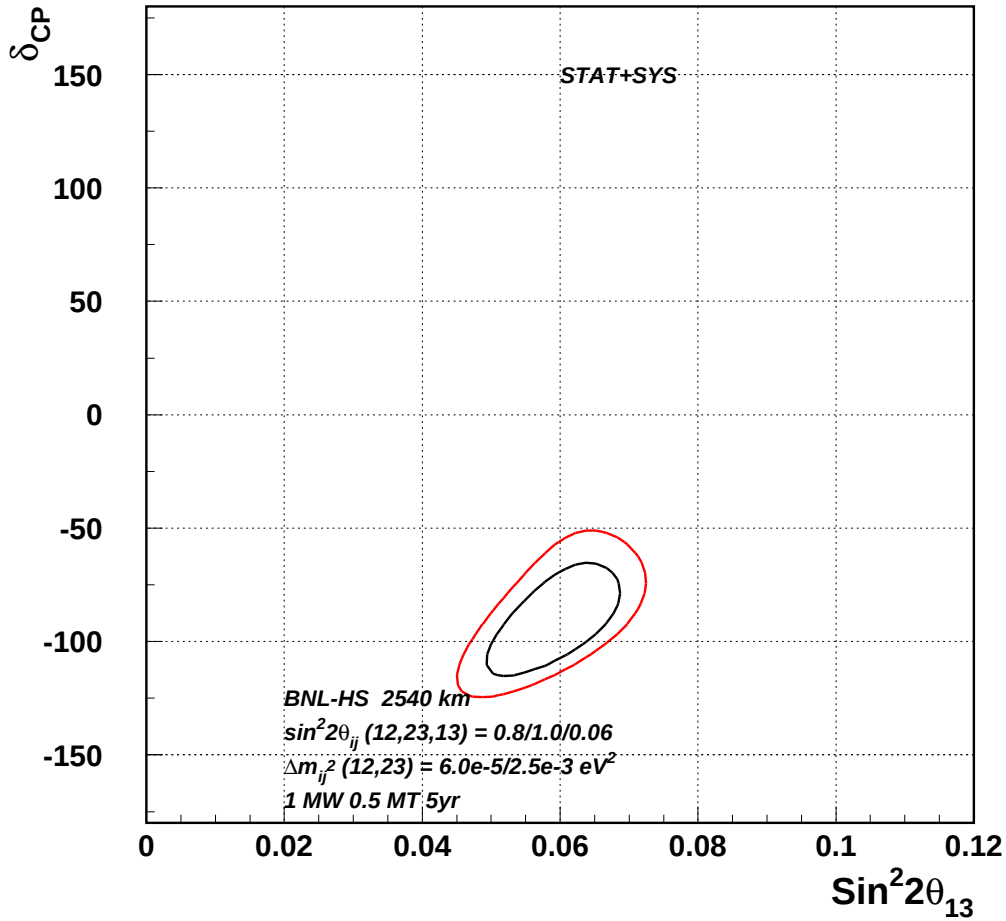
$$\sin^2 2\theta_{12} = 0.8, \sin^2 2\theta_{23} = 1.0$$

$$\delta_{CP} = 45^\circ, \sin^2 2\theta_{13} = 0.06$$

68%, and 90% C.L.

Measurement of $\delta_{CP} = -90^\circ$

90 % C.L. for δ_{CP} vs $\sin^2 2\theta_{13}$



Systematic error of 10% on backg.

$$\Delta m_{21}^2 = 6 \times 10^{-5} \text{ eV}^2, \Delta m_{31}^2 = 2.5 \times 10^{-3} \text{ eV}^2$$

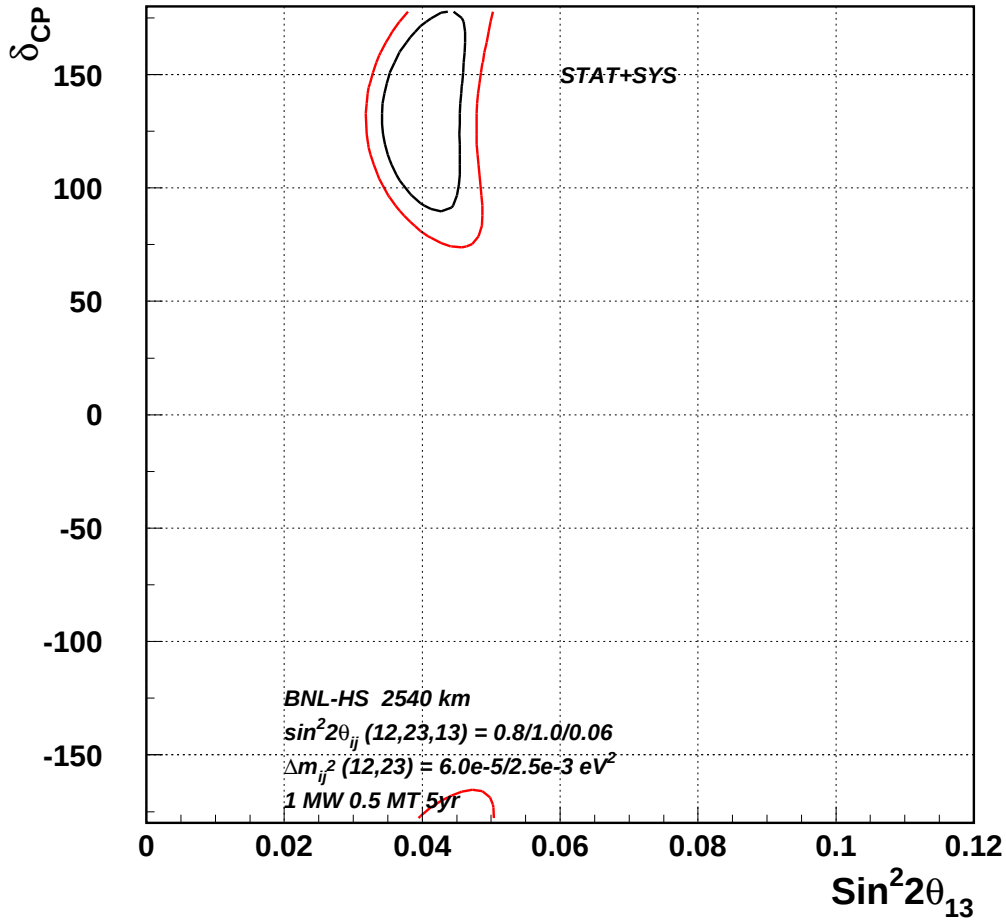
$$\sin^2 2\theta_{12} = 0.8, \sin^2 2\theta_{23} = 1.0$$

$$\delta_{CP} = -90^\circ, \sin^2 2\theta_{13} = 0.06$$

68%, and 90% C.L.

Measurement of $\delta_{CP} = 135^\circ$

90 % C.L. for δ_{CP} vs $\sin^2 2\theta_{13}$



Systematic error of 10% on backg.

$$\Delta m_{21}^2 = 6 \times 10^{-5} \text{ eV}^2, \Delta m_{31}^2 = 2.5 \times 10^{-3} \text{ eV}^2$$

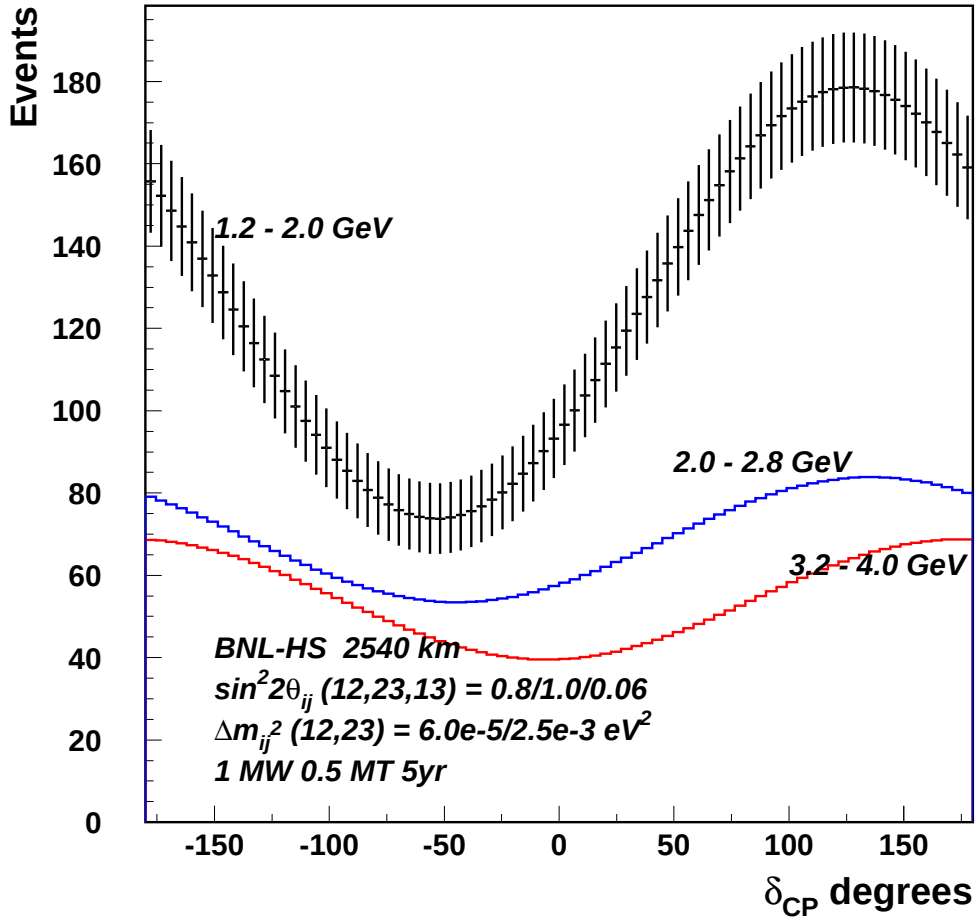
$$\sin^2 2\theta_{12} = 0.8, \sin^2 2\theta_{23} = 1.0$$

$$\delta_{CP} = 135^\circ, \sin^2 2\theta_{13} = 0.06$$

68%, and 90% C.L.

Effect of δ_{CP} on the spectrum.

Effect of δ_{CP} in 3 energy bins



Event rate in 3 energy bins.

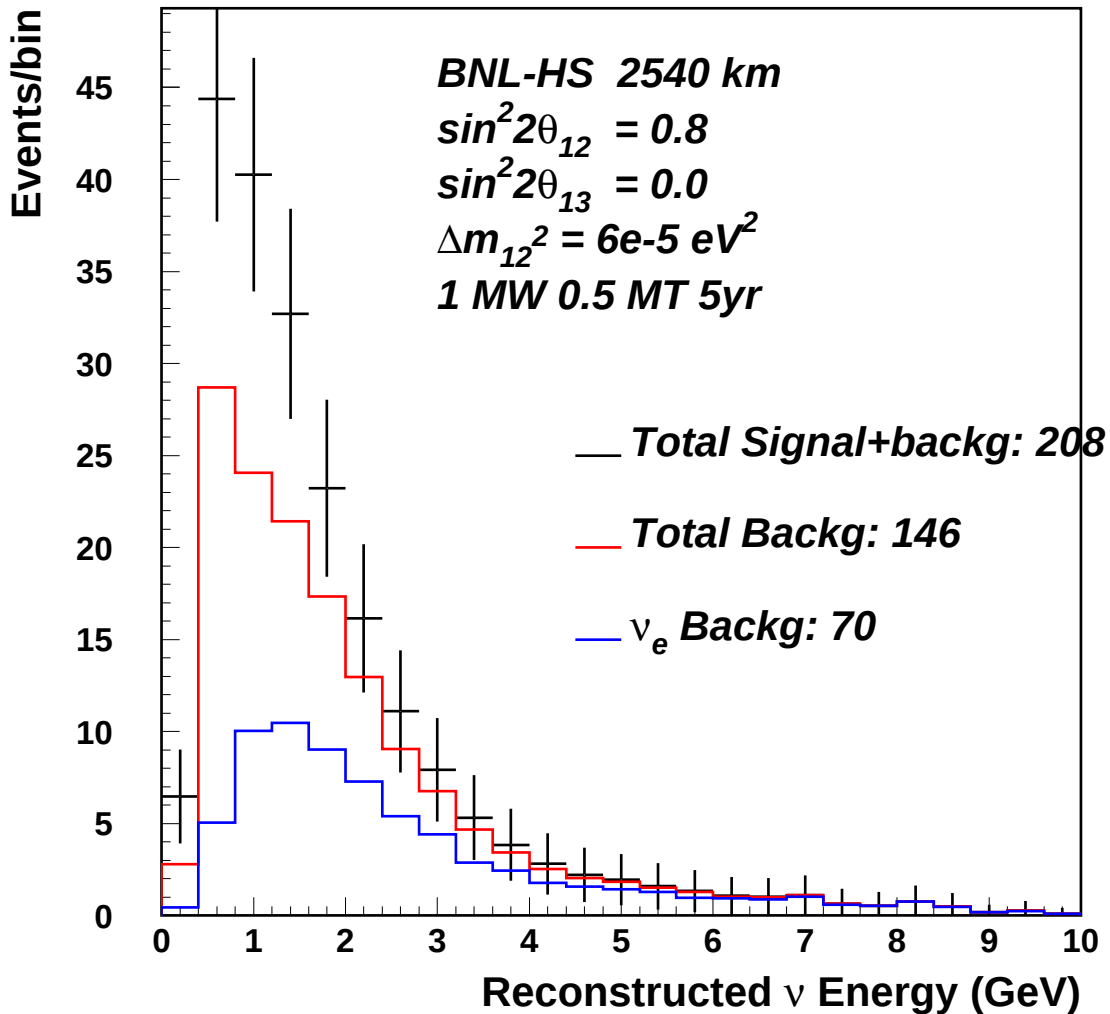
$$\Delta m_{21}^2 = 6 \times 10^{-5} \text{ eV}^2, \Delta m_{31}^2 = 2.5 \times 10^{-3} \text{ eV}^2$$

$$\sin^2 2\theta_{12} = 0.8, \sin^2 2\theta_{23} = 1.0$$

$$\sin^2 2\theta_{13} = 0.06$$

Measurement of Δm_{12}^2

ν_e APPEARANCE FROM Δm_{12}^2 ONLY



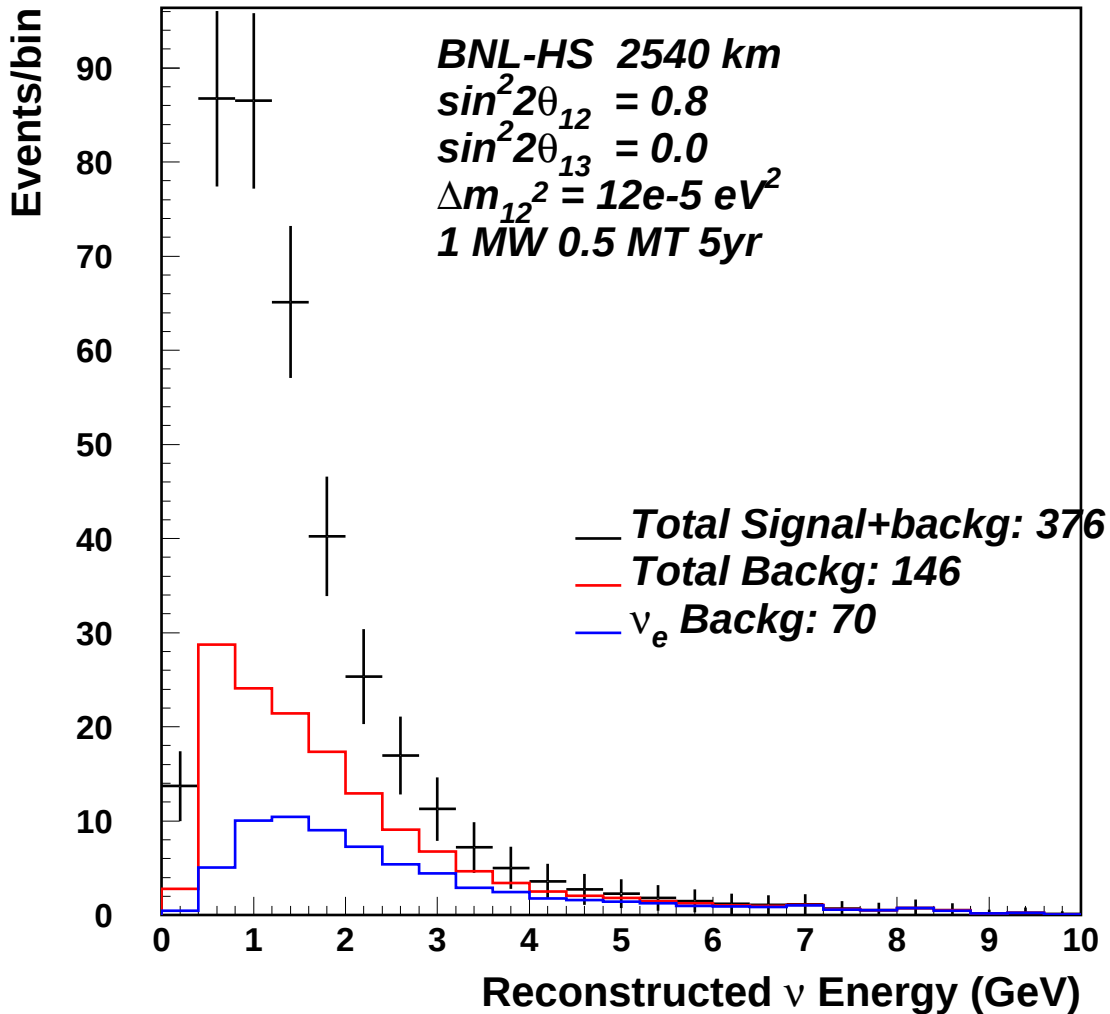
$$\theta_{13} = 0, \Delta m_{12}^2 = 6 \times 10^{-5} \text{ eV}^2$$

Excess of ~ 50 events. Must know background

$$\text{Recall } \sin^2(1.27 \Delta m_{12}^2 2540 \text{ km} / 1 \text{ GeV}) = 0.037$$

Measurement of Δm_{12}^2

ν_e APPEARANCE FROM Δm_{12}^2 ONLY

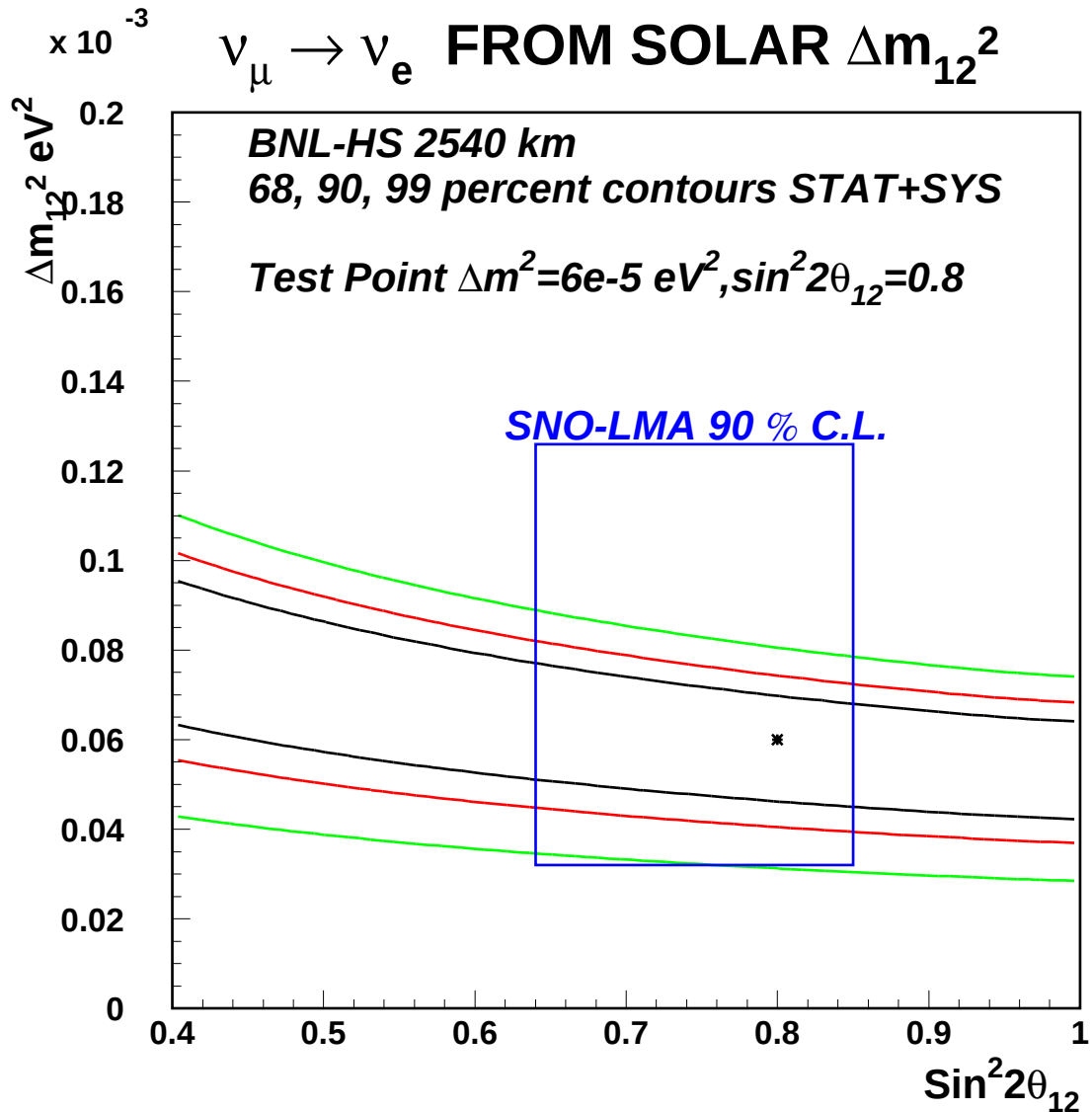


$$\theta_{13} = 0$$

$$\Delta m_{12}^2 = 12 \times 10^{-5} \text{ eV}^2$$

Excess of ~ 230 events. Unmistakable.

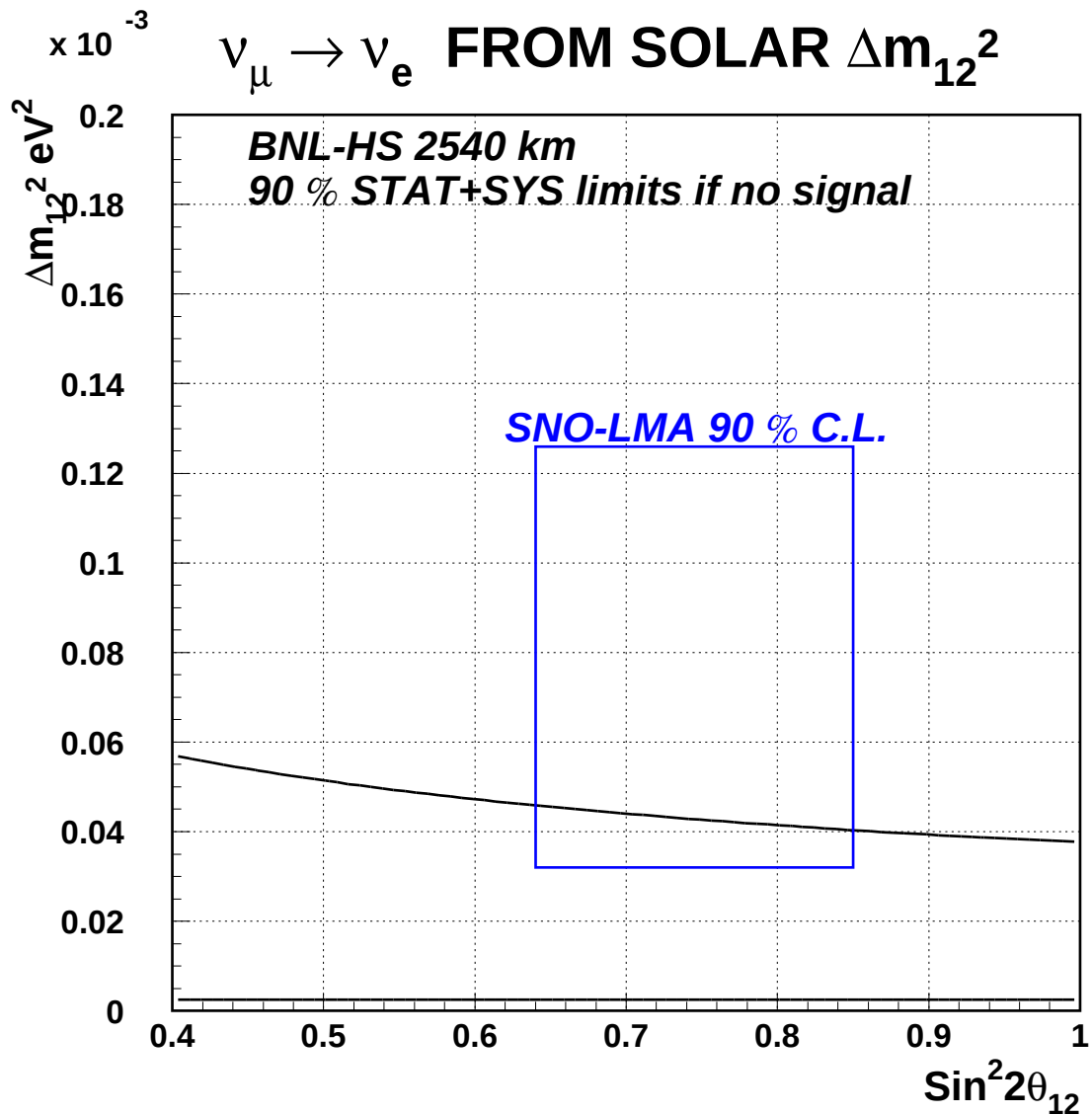
Measurement of Δm_{12}^2



Independent $\sim 15\%$ measurement of Δm_{12}^2

Needs $\sim 10\%$ error on backg. $\Rightarrow$ near detector.

Limit on Δm_{12}^2 vs $\sin^2 2\theta_{12}$



If no signal then a limit can be obtained that almost eliminates LMA.

Analysis Flow Chart

How the experiment will proceed:

- After 2 years of running get a very precise measurement of Δm_{23}^2 from disappearance and definitive signal of oscillations.
- From the measured Δm_{23}^2 predict the shape of the electron spectrum including matter effects.
- Do we have a peak in the electron spectrum at the expected energy ? Yes No
- NO: Either $\sin^2 2\theta_{13}$ too small or inverted mass hierarchy $\Delta m_{32}^2 < 0$.
 - Get an independent measurement of Δm_{12}^2 at about $\pm 15\%$.
 - Run with anti-neutrinos. (next next slide)
- YES: GREAT NEWS ! GOTO NEXT SLIDE.

- YES: There is a peak in the electron spectrum from the neutrino beam.
 - Use Δm_{12}^2 from SNO and KAMLAND and make a fit to the spectrum for CP angle versus $\sin^2 2\theta_{13}$.
 - Accumulate more statistics and make a combined fit for Δm_{12}^2 , δ_{CP} and θ_{13} .
 - Is the CP angle too small ? NO YES
- NO: Finished ! Still run antineutrinos for more precise δ_{CP} .
- YES: Run anti-neutrinos for more sensitivity on δ_{CP} .

Measure both $\sin^2 2\theta_{13}$ and δ_{CP}

- Running with anti-neutrinos if no peak in the electron spectrum from neutrinos

Is there a peak in the electron spectrum from anti-neutrinos ? Yes No

- Yes The mass hierarchy is inverted. Proceed to measure $\sin^2 2\theta_{13}$ and CP angle with anti-neutrinos.
- No $\sin^2 2\theta_{13}$ is too small. Proceed to social work.

If inverted hierarchy;

measure both $\sin^2 2\theta_{13}$ and

δ_{CP} .

OR $\sin^2 2\theta_{13}$ is just too small for conventional beam.

Summary of our study

- Baseline of > 2000 km with wide band conventional beams are the next step in accelerator neutrino physics.
- Extraordinary, large physical effects will be seen in such an experiment.
- Very good sensitivity to neutrino properties.
 - $< 1\%$ resolution on Δm_{32}^2
 - $< 1\%$ resolution on $\sin^2 2\theta_{23}$
 - Sensitivity to $\sin^2 2\theta_{13} \sim 0.005$ over a wide range of Δm_{32}^2
 - Sensitivity to CP parameter $\pm 25^\circ$ with neutrinos alone.
 - Sign of Δm_{32}^2 over a wide range.
 - Measurement of Δm_{12}^2 at $\pm 15\%$
- The electron spectrum has a lot of physics. It can be extracted using some outside information on parameter.

Measurement matrix

Neutrino running only; Running: 5×10^7 sec.

Baseline: 2540 km; beam: 1 MW at 28 GeV; detector: 500 kT

	Δm_{32}^2	$\sin^2 2\theta_{23}$	Δm_{12}^2	$\sin^2 2\theta_{13}$ 90 % C.L.	δ_{CP}
$\Delta m_{32}^2 > 0.001$	$< 1\%$	$\sim 1\%$	$\pm 15\%$	~ 0.005	
$\Delta m_{32}^2 > 0.001$ $\sin^2 2\theta_{13} > 0.01$	$< 1\%$	$\sim 1\%$	$\pm 15\%$	± 0.01	$\pm 25^\circ$
$\Delta m_{32}^2 > 0.001$ $\sin^2 2\theta_{13} < 0.01$	$< 1\%$	$\sim 1\%$	$\pm 15\%$	No Measure.	No Measure.

Not complete story, but an impression. Assume $m_3 > m_2 > m_1$.

Need good energy calibration for Δm_{32}^2 ($\sim 100 MeV$ LINAC ?)

Need small error on backg. for Δm_{12}^2 and CP. (Near Detector)

What is Next ?

Study will be published shortly.

Can we use events such as $\nu_e + N \rightarrow e^- + \pi^+ + N$

Anti-neutrino sensitivity. Hierarchy determination.

Parameter correlations.

Background determination with near det.

- The experiment is technically feasible.

Direct costs.

AGS upgrade, Hill, Proton transp., horns,
decay tunnel: $\sim \$150M$

Detector: \$300 M for 10% PMT coverage.

This can be a staged program that starts
with \$90 M at the AGS and \$150 M at
Homestake for first critical results.

- The detector has applications far beyond
accelerator neutrinos. And should have a very
diverse and rich physics program.