Quantum Logic Gates for Capacitively Coupled Josephson Junctions

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Superconductivity Research





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Outline

- Motivation
- Energy Spectrum
- Logic Gates
- Experimental Prospects
- Conclusion

Theoretical Motivation:

- Current-biased Josephson junctions are seemingly **macroscopic** objects obeying quantum mechanics.
- Fabrication and interactions allow for directed design of **scalable** circuits.
- Electrical properties are easily **measured**.

Experimental Motivation:

- Coherence: Martinis *et al.*, Yu *et al.* (2002) demonstrated Rabi oscillations in single qubit dynamics.
- **Coupling**: Berkley *et al.* (2003) demonstrated tunable coupling through spectroscopy.
- Experimental **control** of coherent two-qubit dynamics of current-biased junctions should be possible—quantum logic gate design is necessary.

Current-Biased Josephson Junction



Capacitively Coupled Josephson Junctions



$$E_{C} = \frac{e^{2}}{2C_{J}}, E_{J} = \frac{\hbar I_{C}}{2e}$$
$$J_{1} = I_{1} / I_{C}, J_{2} = I_{2} / I_{C}$$

$$\zeta = C / (C + C_J)$$

$$H = 4E_{C}(1+\zeta)^{-1}(p_{1}^{2}+p_{2}^{2}+2\zeta p_{1}p_{2})$$
$$-E_{J}(\cos \gamma_{1}+J_{1}\gamma_{1}+\cos \gamma_{2}+J_{2}\gamma_{2})$$

Energy Spectrum

$$N_s = 4, \ \zeta = 0.01$$

 $\sqrt{1 - J_{1,2}} = \sqrt{1 - J_0} (1 \pm \varepsilon)$
 $\hbar \omega_0 = (8E_C E_J)^{1/2} (1 - J_0^2)^{1/4}$

•Energy states are unentangled away from avoided level crossings.

•Entanglement is maximized at the avoided level crossings.



Gate Design

• **Control**: Interactions controllable (tuned on and off) through bias currents for small coupling.

(e.g. $\zeta = 0.01$)

- **Dynamical conditions**: Characteristic ramp time must satisfy $\frac{2\pi}{\omega_0} < \tau_R < \frac{1}{\zeta} \frac{2\pi}{\omega_0} \approx 100 \times \frac{2\pi}{\omega_0}$
- Leakage: Both tunneling and evolution into the auxiliary states |02> and |20> must be taken into account.

$$N_S \ge 4$$

Gate Operation

- Start from detuned junctions.
- Ramp bias currents, in time τ_R , from ε_A to ε_B .
- Wait for time τ_{I} .
- Detune the junctions.



Phase Gate Operation



This avoided level crossing is isolated, so the other two-qubit states $|00\rangle$, $|01\rangle$ and $|10\rangle$ are unaffected.

Phase Gate Simulation











Swap-Like Operation









Fidelity and Leakage

Fidelity is the probability that the gate operation is successful.

Leakage is the probability that the gate drives transitions out of the two-qubit basis.

	Fidelity	Leakage	Time(ns)
Phase	0.996	0.003	14.85
Swap	0.972	0.006	10.7

 $\omega_0/2\pi = 6 \text{ GHz}$

Experimental Demonstration

- The nanosecond pulse times can be generated with conventional electronics.
- Energy levels can be determined spectroscopically.
- Coherence requires characteristic ("RC") dissipation times ≥ 1 µs:
 - C=6pF \Rightarrow R \geq 160 k Ω at GHz frequencies.
 - BUT, bias lines typically have $R \approx 50 \Omega$!

Achieving Coherence: Impedance Transformers





Resistor (Gubrud 2001)

001) LC isolator (Berkley 2002)



Junction (inductor) (Martinis 2002)



Inductor (tunable T₁) (Martinis 2003)

Properties and Optimizations

- Fixed coupling requires undesirably large detuning of junctions.
 - Non-identical qubit energy levels.
 - Possibly large tunneling rates.
- Swap gate uses delicate energy level structure.
 - Sensitive to bias current noise.
 - Difficult with non-identical junctions.
- Gate times are close to time scales for high-fidelity single-qubit operations.

Three-Junction Scheme



•Junction B is used to entangle junctions A₁ and A₂.

•Before and after the operation, junction B is in its ground state.



Cf. Blais, Maassen van den Brink, Zagoskin PRL (2003)

Conclusion

- Designed and numerically simulated two fundamental quantum logic operations, each with fidelity F>0.97.
- Explored multiple junction schemes for controlled coupling.
- Experimental demonstration of these logic gates is possible.

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