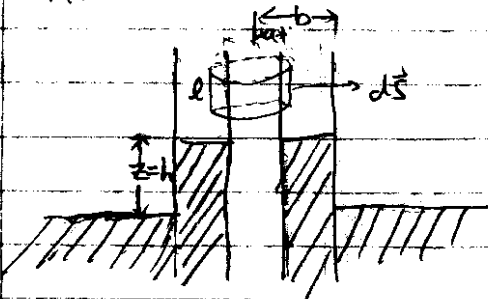


## HW#8 solution.

12.



By symmetry,  $\vec{E}$  should be independent of azimuthal angle and along radial direction.

From Gauss theorem, in the vacuum

$$\oint \vec{E} \cdot d\vec{s} = \frac{Q}{\epsilon_0}$$

$$\Rightarrow E 2\pi r l = \frac{Q}{\epsilon_0}$$

$$\Rightarrow E = \frac{\lambda}{2\pi r \epsilon_0}, \quad \lambda = \frac{Q}{l}$$

Since the voltage across the two conductor is  $V$ ,

$$\int_a^b \vec{E} \cdot d\vec{r} = \int_a^b \frac{\lambda}{2\pi r \epsilon_0} dr = V$$

$$\Rightarrow \lambda = \frac{2\pi \epsilon_0 V}{\ln\left(\frac{b}{a}\right)}$$

$$\Rightarrow E = \frac{V}{r \ln\left(\frac{b}{a}\right)}$$

Clearly the electric field inside the medium is the same since  $\epsilon$  will be cancelled.

So, the electric energy contributed by the medium inside the conductor is

$$\begin{aligned} W_e &= -\frac{1}{2} \int dV (\epsilon - \epsilon_0) E^2 \\ &= -\frac{1}{2} \int_0^z \int_0^{2\pi} \int_a^b (\epsilon - \epsilon_0) \left[ \frac{V}{r \ln\left(\frac{b}{a}\right)} \right]^2 r dr d\phi dz \\ &= -\frac{V^2}{\ln\left(\frac{b}{a}\right)} \pi (\epsilon - \epsilon_0) z \quad \text{①} \end{aligned}$$

where  $z$  is the height of the medium inside the conductor.

If the system reaches equilibrium, the electric force will be equal to the gravity of the medium, i.e.

$$-\frac{\partial W_e}{\partial z} = \rho g \pi (b^2 - a^2) h \quad \textcircled{2}$$

By ①, ②, one get

$$\chi_e = \frac{\epsilon - \epsilon_0}{\epsilon_0} = \frac{\rho g (b^2 - a^2) h \ln \frac{b}{a}}{\epsilon_0 V^2}$$

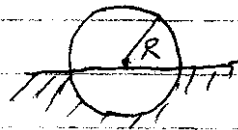
2) stress for fields perpendicular to a surface is

$$T = \frac{\vec{D} \cdot \vec{E}}{2}$$

$$\vec{T} = \frac{\vec{D} \cdot \vec{E}}{2}$$

From HW #7 (Jackson 4.10)

$$\vec{E} = \frac{Q}{2\pi(\epsilon + \epsilon_0)r^2} \hat{r} \quad \textcircled{1}$$

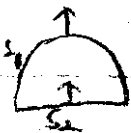


If the system reaches equilibrium

$$\oint \vec{T} \cdot d\vec{S} + mg \hat{z} + \frac{1}{2} \rho g \left( \frac{4}{3} \pi R^3 \right) = 0 \quad mg = \frac{4}{3} \pi R^3 \rho g f_0$$

$$\Rightarrow \int_{\text{up}} \frac{\epsilon_0 E^2}{2} d\vec{S} + \int_{\text{down}} \frac{\epsilon E^2}{2} d\vec{S} = -\rho g \frac{4}{3} \pi R^3 \left( \frac{1}{2} - f_0 \right)$$

$$\Rightarrow \frac{1}{2} E^2 \left[ \int_{\text{up}} \epsilon_0 d\vec{S} + \int_{\text{down}} \epsilon d\vec{S} \right] = -\rho g \frac{4}{3} \pi R^3 \left( \frac{1}{2} - f_0 \right) \quad \textcircled{2}$$



consider a closed surface as shown to the left

$$\oint \phi d\vec{S} = \int \nabla \phi \cdot d\vec{V}$$

$$\Rightarrow \oint d\vec{S} = 0 \Rightarrow \int_{S_1} d\vec{S} = \int_{S_2} d\vec{S}$$

So, ② can be rewrite as

$$\frac{1}{2} E^2 (\epsilon_0 - \epsilon) \pi R^2 = -\rho g \frac{4}{3} \pi R^3 \left( \frac{1}{2} - f_0 \right) \quad \textcircled{3}$$

plug ① into ③, one get

$$Q = 2\pi(\epsilon + \epsilon_0) R^2 \left[ \frac{8}{3} \rho g R \frac{\frac{1}{2} - f_0}{\epsilon - \epsilon_0} \right]^{\frac{1}{2}}$$