

Encounters with Supergavity

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Plan : Something new & Something Old

- OLD {
- 1) Intro : Why I failed to invent SUGRA
 - 2) Killing Vectors & Killing Spinors
 - 3) de Sitter, Anti de Sitter & whatever happened to Twistor theory
- NEW {
- 4) Special Holonomy, Bries, Calibrations & Cohomogeneity one metrics

INFINITESIMAL SUPERSYMMETRY

(Killing Spinor)

v

INFINITESIMAL ISOMETRY

(Killing Vector)

$$\delta_K g_{\mu\nu} = \frac{1}{K} g_{\mu\nu} = 0 \Leftrightarrow \boxed{K_{\mu;\nu} + K_{\nu;\mu} = 0}$$

Killing's Eqns.

$$\delta_\epsilon \psi_\mu = 0 \Leftrightarrow$$

$$\boxed{\hat{\nabla}_\mu \epsilon = 0}$$

$$\boxed{\hat{\nabla}_\mu \epsilon = \nabla_\mu \epsilon + E_\mu \epsilon}$$

SUSY ALG \Rightarrow

$$\boxed{K^\mu = \bar{\epsilon} \gamma^\mu \epsilon \quad \text{is Killing vector}}$$

Dirac's Lemma : κ^μ is everywhere future directed
timelike or null.



$$\kappa^0 = \psi^\dagger \psi > 0 \text{ in all frames}$$

\Rightarrow Any BPS solution must be locally stationary
& may admit at most a degenerate horizon



X

$$Q = 0$$



X

$$M > |Q|$$



✓

De Sitter v. Anti-de Sitter



$E^{4,1}$

COSMOLOGICAL EVENT HORIZON

$$-(X^0)^2 + (X^1)^2 + (X^2)^2 + (X^3)^2 + (X^4)^2 = 1$$

X



$E^{3,2}$

GLOBALLY
STATIC

$$-(X^0)^2 + (X^1)^2 + (X^2)^2 + (X^3)^2 - (X^4)^2 = -1$$

✓

Wigner's Problem

$$\forall H \in \mathfrak{L}(\mathfrak{SO}(4,1)) \quad \exists g \in \mathfrak{SO}_0(4,1) \text{ s.t.}$$

$$-H = g H g^{-1}$$

\Rightarrow no positive energy \Rightarrow no susy

$$X^0 = \sqrt{1-r^2} \sinh t$$

$$X^4 = \sqrt{1-r^2} \cosh t$$

$$X^i = r n^i, \quad n^i \in S^2$$

$$ds^2 = -(1-r^2) dt^2 + \frac{dr^2}{1-r^2} + r^2 (d\theta^2 + \sin^2\theta d\phi^2)$$

$$L_{04} = \frac{\partial}{\partial t}$$



$$g = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 & -1 \end{pmatrix} \begin{matrix} \in \mathfrak{SO}_0(4) \\ \in \mathfrak{SO}_0(4,1) \end{matrix}$$

g interchanges inside & outside of cosmological horizon.

Actually $\text{SUGRA} \stackrel{?}{\Rightarrow}$ STRONG ENERGY CONDITION
WHICH EXCLUDES ACCELERATION

Twistor Eqn

$$(\nabla_\mu \gamma_\nu + \nabla_\nu \gamma_\mu) \epsilon = \frac{1}{n} g_{\mu\nu} \gamma^\mu \nabla_\mu \epsilon$$

$$\nabla_\mu \epsilon = \frac{1}{2a} \gamma_\mu \epsilon$$

$$\frac{1}{a^2} + \frac{\Lambda}{3} = 0$$

$\Lambda < 0$; a real : $SO(3, 2) \subset SO(4, 2)$

$\Lambda > 0$; a pure imag : $SO(4, 1) \subset SO(4, 2)$

$\Lambda < 0 \quad \rightarrow \quad \epsilon \gamma^\mu \epsilon \quad \text{Killing Vector}$

$\Lambda > 0 \quad \quad \quad \bar{\epsilon} \gamma^\mu \epsilon \quad \text{Conformal Killing Vector}$

Twistor Theory is contained within
Conformal Supergravity.

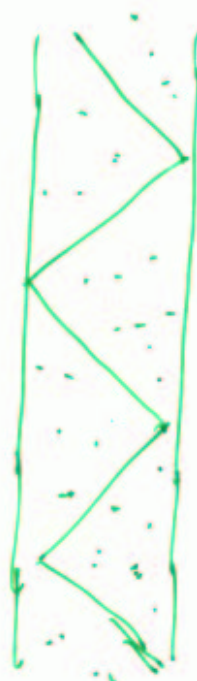
D3 - Branes or (M2 or M3)

$$ds^2 = -H^{-1/2} (-dt^2 + dx_3^2) + H dy_6^2$$

$$\nabla_y^2 H = 0$$

$$F_5 = d \left(\frac{dt dx_1 dx_2 dx_3}{H} \right) + * (d(\dots))$$

$$H = 1 + \frac{1}{|y|^4}$$



No horizons formed
in $AdS_5 \times S^5$

Spatial interpolation between
2 maximal susy vacua
 $\frac{1}{2}$ BPS.

M5 is like D3 non singular

M2 is like Reissner-Nordstrom : singular sources

Special Holonomy

Killing spinors of D3 are constructed from constant or parallel spinors of \mathbb{E}^6

(M2 $\rightarrow \mathbb{E}^9$, M5 $\rightarrow \mathbb{E}^5$ etc..)

Idea 1): replace \mathbb{E}^n by curved mfd

M^n admitting constant spinors

$$\nabla \epsilon = 0.$$

(ii) add additional NS & RR fluxes

$$d * F = i F_3 \wedge \bar{F}_3$$

$$\nabla^2 H = |F_3|^2$$

(Fractional branes)

need to find (M, g) & harmonic forms

Programme of wk. with Cvetič, Liu & Pope)

Special & Exceptional Holonomy

Bergu. 1955 M irreducible & not symmetric

$\text{Hol}(M, g)$

$U(\frac{n}{2}) \subset SO(n)$: Kähler

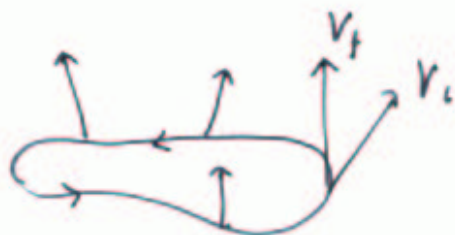
$SU(\frac{n}{2}) \subset SO(n)$: Calabi-Yau

$Sp(\frac{n}{4}) \subset SO(n)$: HyperKähler

$G_2 \subset SO(7)$: Exceptional

$Spin(7) \subset SO(8)$

$Sp(\frac{n}{4}) \cdot Sp(1) \subset SO(n)$: Quaternionic



$$V_t = R V_f$$

$$R \subset \text{Hol} \subseteq SO(n)$$

Group generated by parallel transport:

$$de + \theta \wedge e = 0$$

$$R = d\theta + \theta \wedge \theta$$

: Lie algebra valued 2-form
value in $\text{Hol}(M, g)$

Holonomy gps which present one or more spinors

Susy's

$SU(\frac{n}{2})$: 2 constant spinors

$N=2$

$Sp(\frac{n}{4})$: $n+1$

$N=n+1$

G_2 : 1

$N=1$

$Spin(7)$: 1 : ..

$N=1$

i) $\nabla_{\mu} \epsilon = 0 \Rightarrow Ric = 0$
[not true for Lorentz signature!]

ii) Hol is subgroup of $SO(n)$ leaving invariant a spinor

Eg. spinor rep of $SO(7)$: $\underline{8} = \underline{1} + \underline{7}$

↑ preserved by $G_2 \subset SO(7)$

where in fact G_2 preserves the structure constants of the octonions.

Calibrating Forms

$$\bar{e} \gamma_{\mu_1 \dots \mu_p} \in = \bar{\Phi}_{\mu_1 \dots \mu_p} = \frac{1}{p!} \bar{\Phi}_{\mu_1 \dots \mu_p} dx^{\mu_1} \wedge \dots \wedge dx^{\mu_p}$$

$$\nabla \bar{e} = 0 \Rightarrow \nabla \bar{\Phi} = 0 \Rightarrow d\bar{\Phi} = 0.$$

in fact $\bar{\Phi}$ [tangent plane] \leq |tangent plane|

now  $\int_S \bar{\Phi} = \int_{S'} \bar{\Phi}$

S is calibrated if $\bar{\Phi}|_S = \text{volume form}$

$$\Rightarrow \int_S \bar{\Phi} = \text{vol}(S) \leq \text{vol}(S')$$

S minimizes volume in its homology class.

SUSY cycles & Calibrating forms

BI actions have K -symmetry

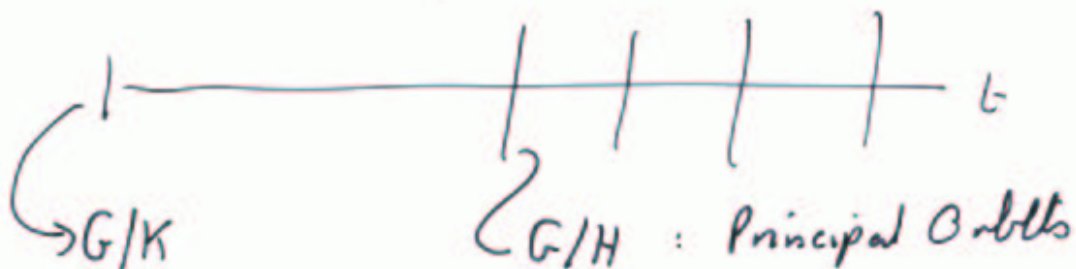
$$K \text{ symmetry} \Rightarrow \text{SUSY} \Rightarrow \text{Calibrated}$$

$SU(\frac{n}{2})$	ω \leftarrow Kähler Ω \leftarrow holomorphic 3-form \Rightarrow <u>holomorphic cycles</u> <u>or SLAGS</u> ; $\omega _S = 0$; $\Re \Omega _S = \text{vol}$ SLAGS have dimension $\frac{n}{2}$
G_2	Φ \leftarrow associative 3-form $*\Phi$ <u>associative 3-fold</u> <u>co-associative 4-fold</u>
$Spin(7)$	$\Psi = *\Psi$ self-dual 4-form <u>Cayley man 4-fold</u>

Cohomogeneity one examples

G acts by isometries with co-dimension one principal orbits.

$$M = G/K \sqcup (\mathbb{R}_+ \times G/H)$$

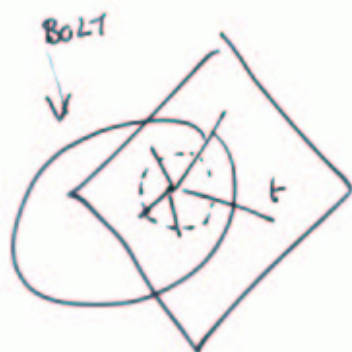


degenerate orbit
(BOLT)

$$K \supset H$$

$$K/H = S^d$$

Globally M^n is an \mathbb{R}^{d+1} bundle over G/K
 \uparrow zero section



Example $M = T^* S^p$

$\dim M = 2p$

$$G = SO(p+2)$$

$$K = SO(p+1)$$

$$H = SO(p)$$

$p=3$ is "deformed" coset

Stenzel metrics

Einstein Eqns

$$ds^2 = dt^2 + g_{ab}(t) e^a(x) e^b(x)$$

$a=1 \dots \dim G/H$

e^a G-covariant vielbein

Cartan-Maurer
Forms

$$\mathfrak{g} = \mathfrak{k} \oplus \mathfrak{p}$$

$\{e^a\}$ are a basis for \mathfrak{p}

$g_{ab}(t)$ is a curve of G-invariant metrics on G/H



$$L = T - V \quad ; \quad V$$

\uparrow Ricci scalar of G/H

$$G_{AB} \frac{dq^A}{dt} \frac{dq^B}{dt} - V(q^A)$$

\uparrow DeWitt metric (Lorentzian signature!)

$$T = T_{ij} g^{-1}{}^i{}_k s^{-1}{}^k{}_j - \frac{1}{2} (T_{ij} s^{-1}{}^i{}_j)^2$$

2nd order ODEs plus one constraint

\Rightarrow

$$\boxed{T + V = 0}$$

DIFFICULT TO
SOLVE

Special Holonomy & Superpotentials

We find that imposing Special holonomy
 \Rightarrow 1st order eqns which themselves imply
the 2nd order eqns.

There are many ways of obtaining 1st order
ode's but in all cases we have found
a superpotential

$$V = - G^{AB} \frac{\partial W}{\partial q^A} \frac{\partial W}{\partial q^B}$$
$$\Rightarrow \frac{dq^A}{dt} = G^{AB} \frac{\partial W}{\partial q^B}$$

In a number of cases there may be solved
explicitly & completely (e.g. Spin(7) 3 funcn
ansatz). In others they must be integrated
numerically.

Behaviour nr. in finitly

We typically encounter 2 types of asymptotes although other possibilities can certainly arise if we drop the cohomogeneity one assumption

(i) AC



Asymptotically Conical

$$ds^2 = dt^2 + t^2 g^\infty(G/H)$$

Ricci flat cone where $g^\infty(G/H)$ is ^{Einstein} metric

with Killing spinors

e.g.

<u>Einstein Sasasaki</u>	5-dim	Calabi-Yau 6
<u>Nearly Kähler</u>	6-dim	Calabi-Yau 7
<u>Weak G₂</u>	7-dim	Spin 7
<u>Tri-Sasasaki</u>	8-dim	HK
<u>Einstein Sasasaki</u>	7-dim	Calabi-Yau 8

In 8-dimensions we can have

$$\begin{aligned} \text{Spin}(7) &\cong \text{Spin}(7) \\ \text{SU}(4) &\cong \text{Spin}(6) \\ \text{Sp}(2) &\cong \text{Spin}(5) \end{aligned}$$

curves

(11) Asymptotically Locally Conical ALC

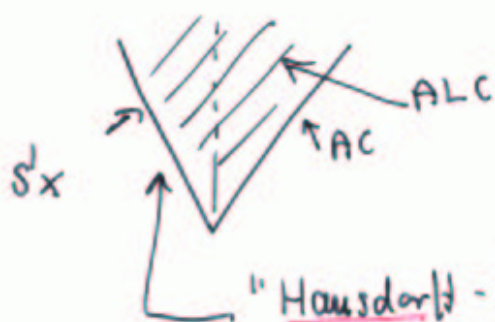
This is more generic & favoured by the eqns.



$$\frac{S^1 \times R \times (G/H \times SO(2))}{AC}$$

a circle of constant length nr. ∞ .

There allows an M-Theory interpretation!



Typical moduli space

"Hausdorff - Cromov limit" in which a circle splits off as a product.

Calibrated BOLTS We typically find

that the BOLT or degenerate orbit is a calibrated submfd.

Example

G-dim C.Y.

cone over $T^{(1,1)} \cong S^1$ bundle over $S^2 \times S^2$



$O(2) \times O(2)$ on CP^1 cplx line bundle.

stenzel: BOLT is SLAG.

Orientifold Planes & D6-Planes

Taub-NUT

v

Atiyah-Hitchin

D6

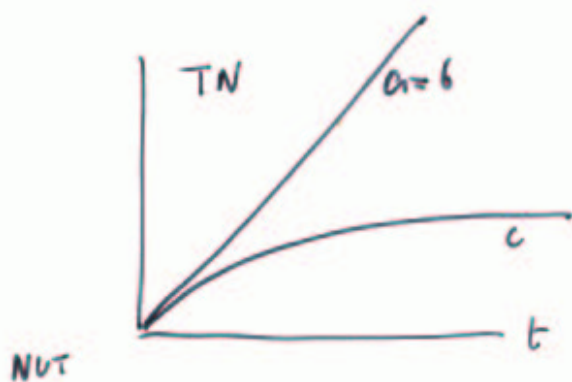
O6

+ve mass

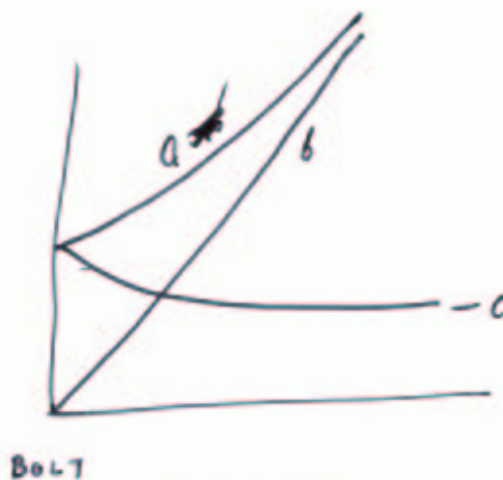
-ve mass

$$ds^2 = -dt^2 + a^2 \sigma_1^2 + b^2 \sigma_2^2 + c^2 \sigma_3^2$$

$$d\sigma_1 = \sigma_2 \wedge \sigma_3 \quad \text{etc.}$$



+ve mass R^4



~~T~~
 $S(S^2)$

Asymptotically AH \sim TN (-ve mass) / CP

$$p: (r, \theta, \phi) \rightarrow (r, \pi - \theta, \phi + \pi)$$

$$c: \psi \rightarrow -\psi$$

\leftarrow M-theory circle.

New Spin 7 examples

Old core
Bryant Salamon

AC example on

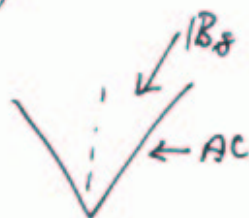
$S^+(S^4)$

↑ chiral spinors

$G = SO(5)$

S^4 : Cayley 4-fold

New examples : i) ALC on $S^+(S^4)$
(explicit) (with parameters)



(ii) isolated ALC core on $\mathbb{R}^8 \Rightarrow A_7$

(iii) ALC example on ; \mathbb{C}_8
bundle over $\mathbb{C}P^2 \leftarrow S^2$

\downarrow_{S^4}

$G/H \cong SO(5)/SO(3) \cong S^7$

wrapped orientifolds